

# Silicon sensors R&D for High energy Physics

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and Medical Imaging

ATLAS PIXEL B LAYER INSERTION



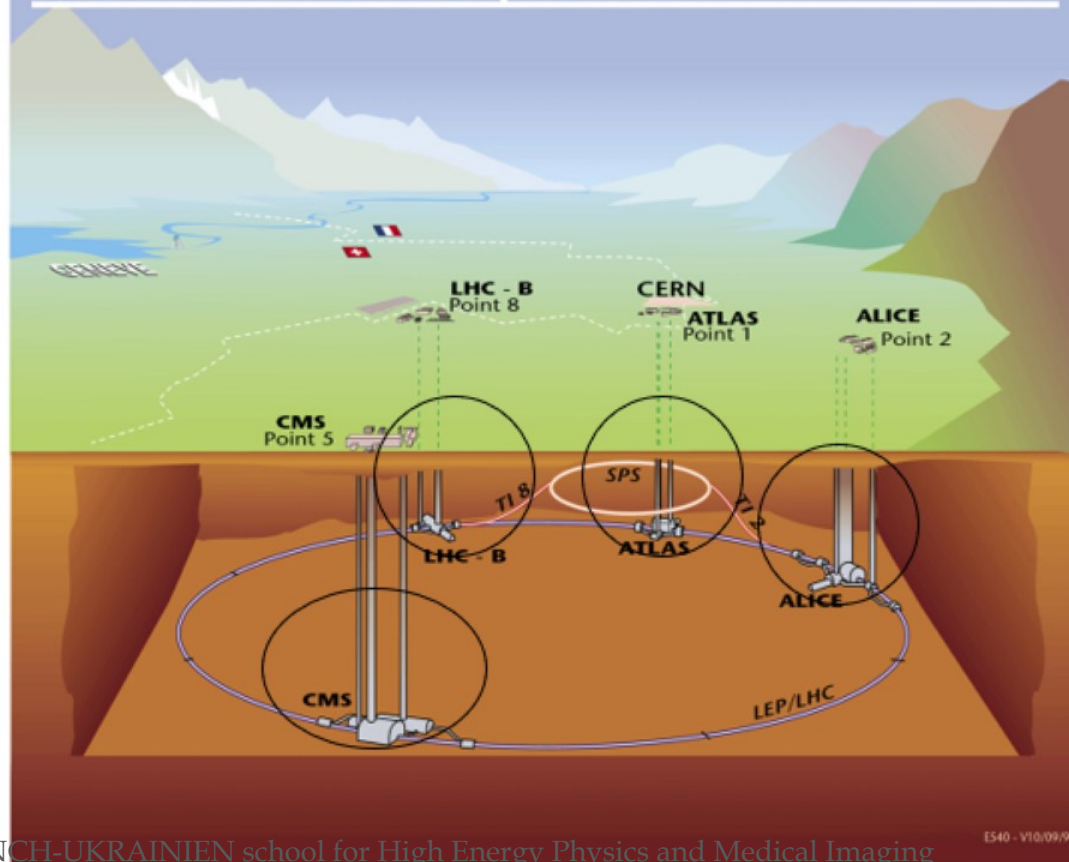
# Plan

- Introduction
- What do we measure with silicon ?
- Basic considerations
- Key parameters
- Silicon pixels for ATLAS
  - Overview of R&D developments
- Conclusions



The Large Hadron Collider (LHC) will be the most powerful instrument ever built to investigate particles properties.

Overall view of the LHC experiments.



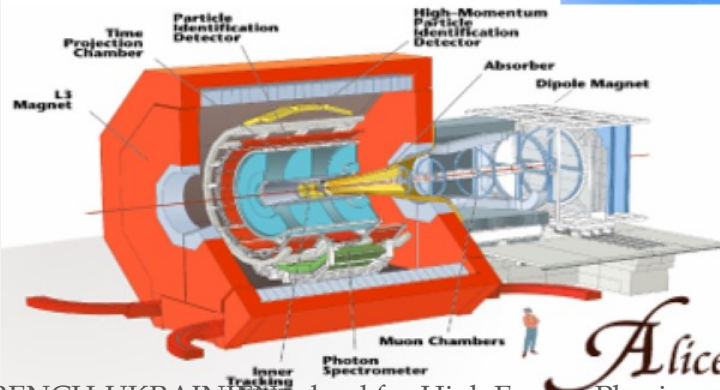
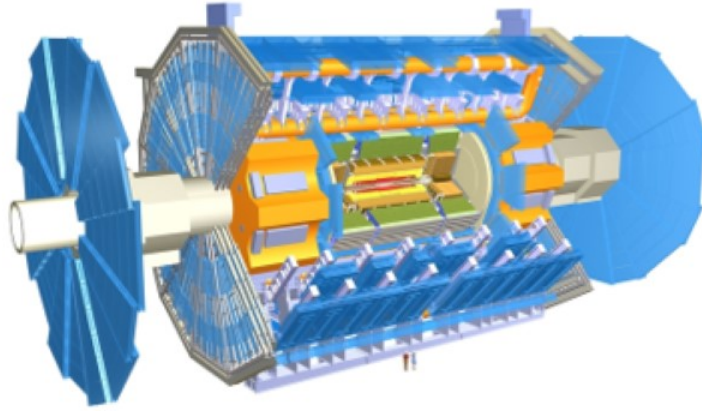
- Four **gigantic underground caverns** to host the huge detectors
- The **highest energy** of any accelerator in the world
- The **most intense beams** of colliding particles
- It will operate at a temperature **colder than outer space**

# The ATLAS Collaboration

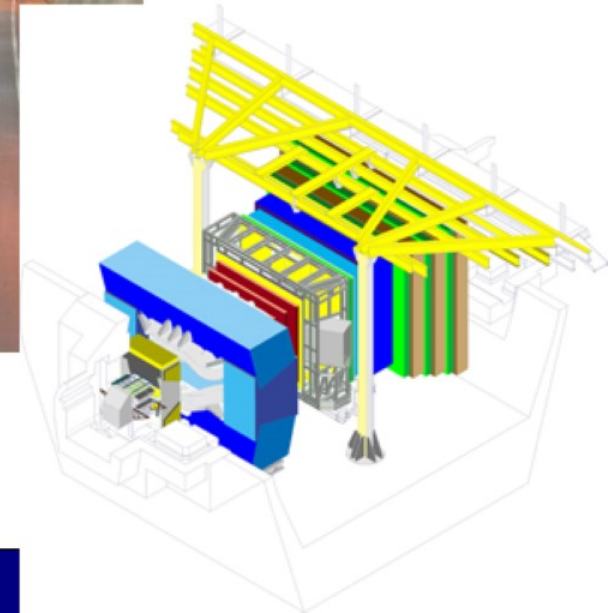
36 Countries  
165 Institutions  
2000 Scientific Authors







*Alice*

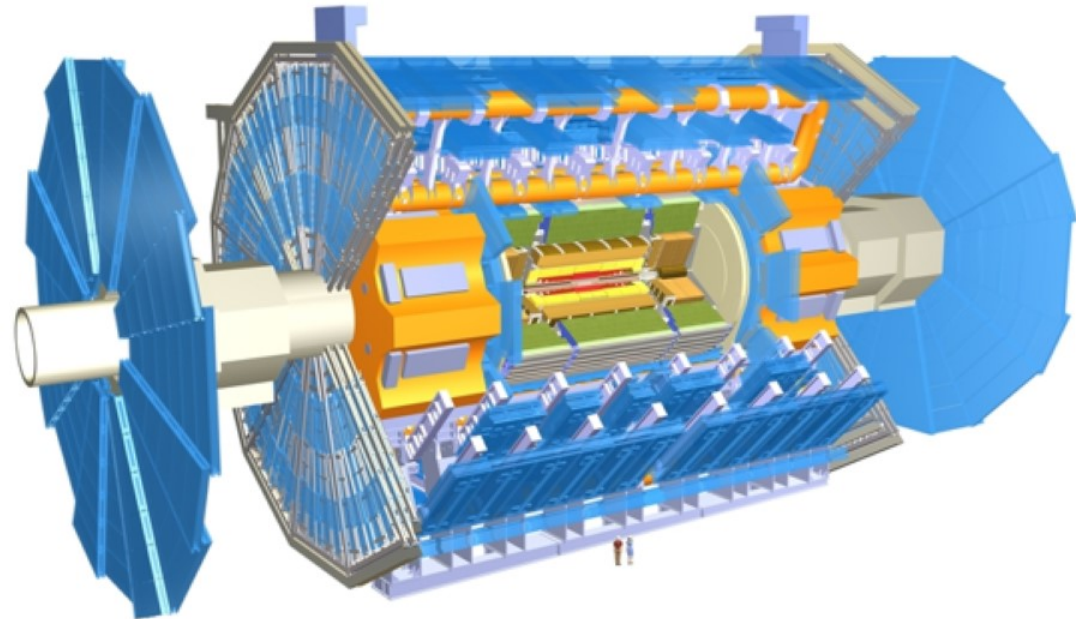




# *The ATLAS Detector*

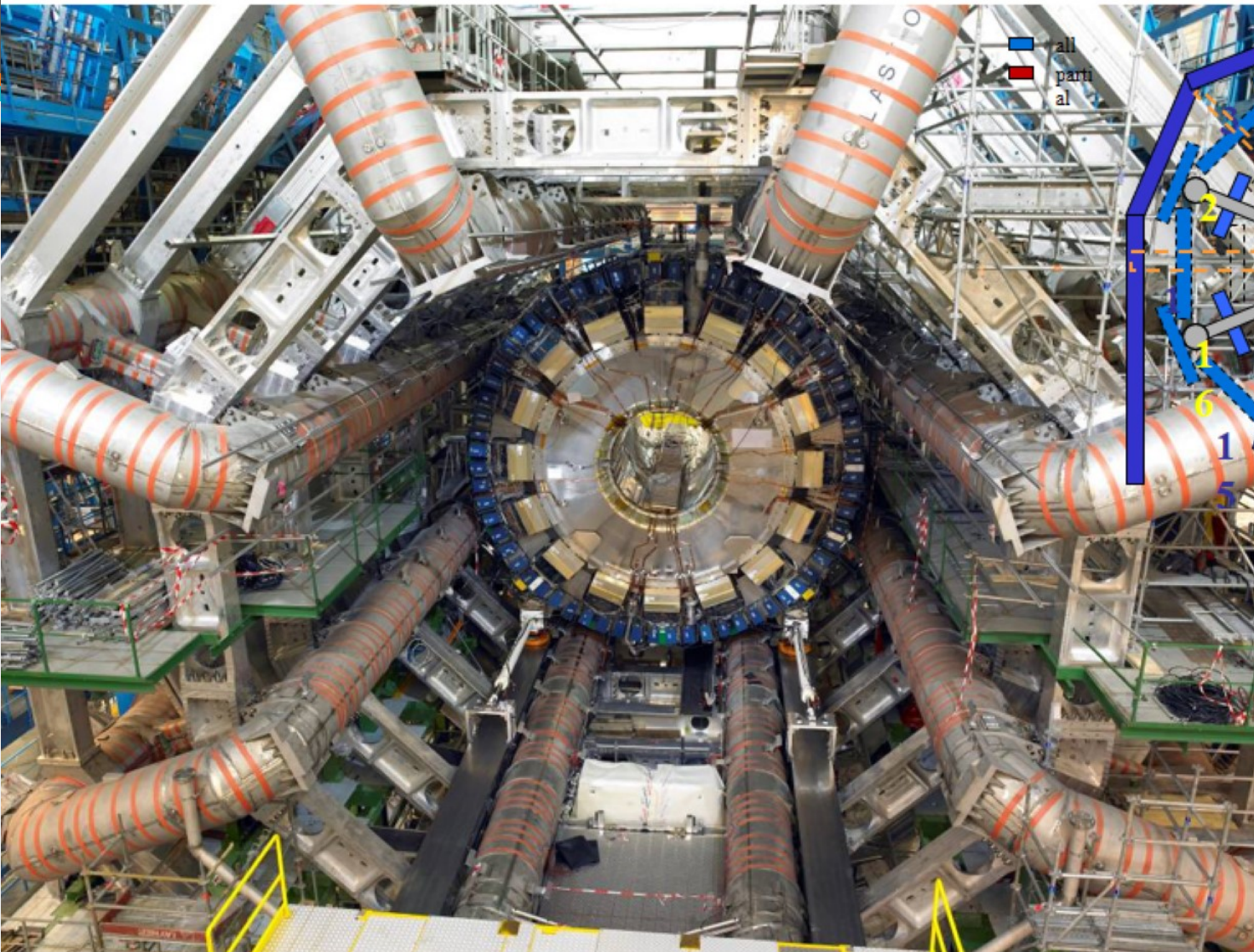


ATLAS superimposed to  
the 5 floors of building 40

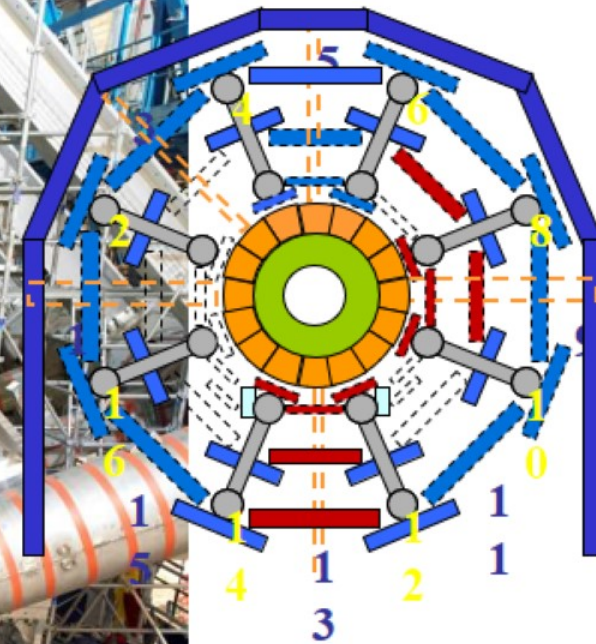


<i>Diameter</i>	<i>25 m</i>
<i>Barrel toroid length</i>	<i>26 m</i>
<i>End-cap end-wall chamber span</i>	<i>46 m</i>
<i>Overall weight</i>	<i>7000 Tons</i>





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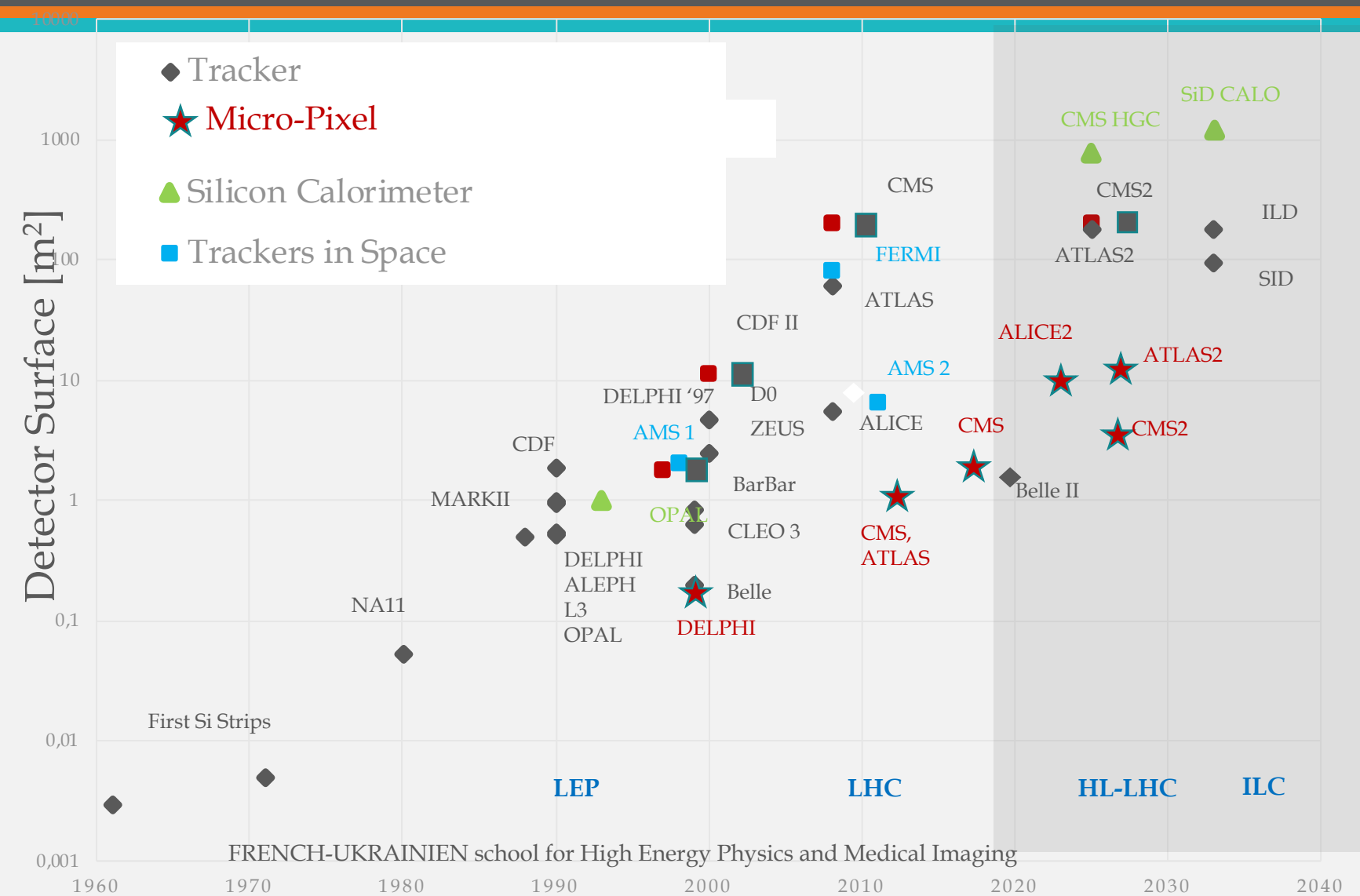
# Inner silicon tracker





# Silicon as material for Trackers! Actual situation

Detector surface  
Globally stays constant  
Pixel area goes up  
significantly  
Cell size goes down significantly



Detector	Strip length [cm]	Strip length [cm]
	pixel size [μm <sup>2</sup> ]	pixel size [μm <sup>2</sup> ]
ATLAS pixel	50x400	25x100 (50x50)
CMS pixel	150x100	50x50 (25x100)
VELO	1 to 7 cm	55x55
ALICE	50X425	28x28

Detector	Channels	Channels
CMS strips	9.8M	42M + 172M
CMS Pixels	127M	2GP
ATLAS strips	6.3M	60M
ATLAS pixels	92M	5GP
VELO	171k	41M
ALICE	12.5M	12.5G

# What is measured

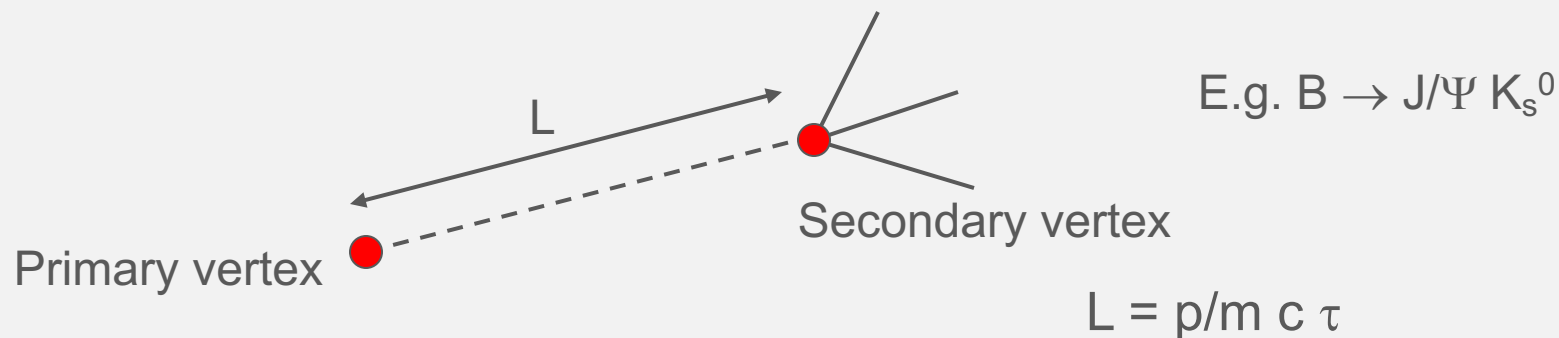
- Measure space points
- Deduce
  - Vertex location
  - Decay lengths
  - Impact parameters

# Signature of Heavy Flavours

Stable particles $\tau > 10^{-6} \text{ s}$		$c\tau$
n		2.66km
$\mu$		658m
Very long lived particles $\tau > 10^{-10} \text{ s}$		
$\pi, K^{\pm}, K_L^0$	$2.6 \times 10^{-8}$	7.8m
$K_S^0, E^{\pm}, \Delta^0$	$2.6 \times 10^{-10}$	7.9cm
Long lived particles $\tau > 10^{-13} \text{ s}$		
$\tau^{\pm}$	$0.3 \times 10^{-12}$	91 $\mu\text{m}$
$B_d^0, B_s^0, \Delta_b$	$1.2 \times 10^{-12}$	350 $\mu\text{m}$
Short lived particles		
$\pi^0, \eta^0$	$8.4 \times 10^{-17}$	0.025 $\mu\text{m}$

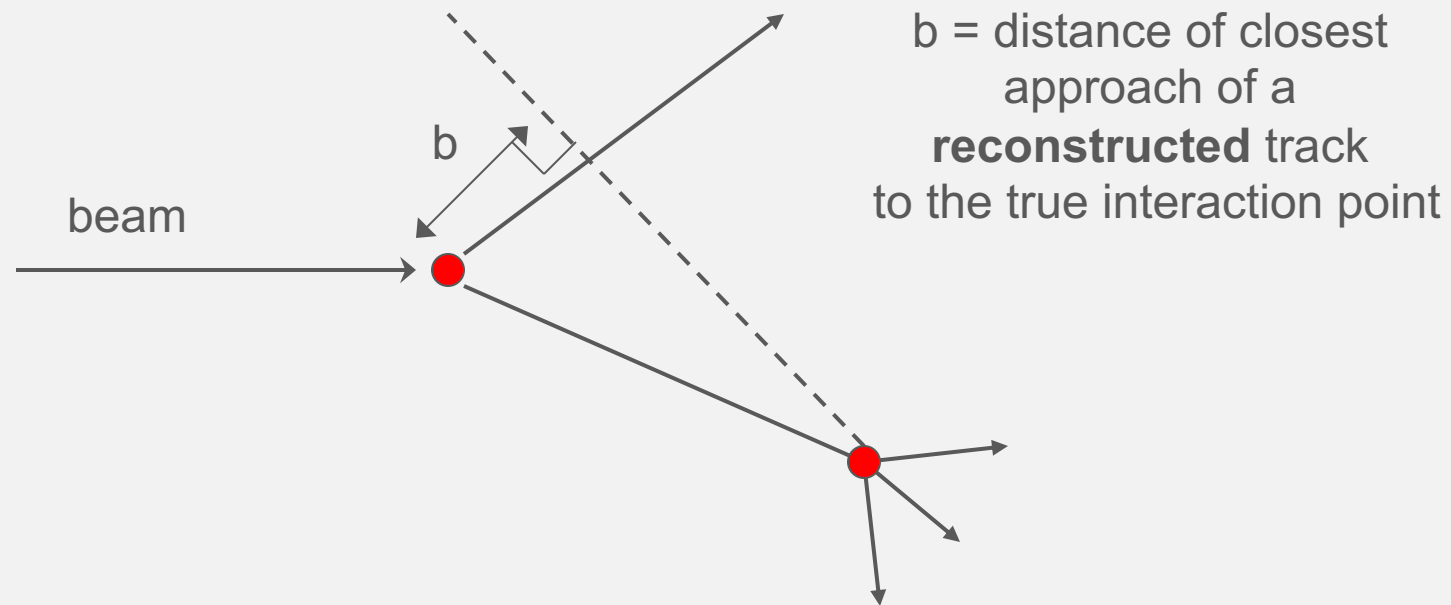


# Decay lengths



- By measuring the decay length,  $L$ , and the momentum,  $p$ , the lifetime of the particle can be determined
- Need accuracy on both production and decay point

# Impact parameter ( $b$ )



# Impact parameter

- Error in impact parameter for 2 precision measurements at  $R_1$  and  $R_2$  measured in two detector planes:

$$\sigma_b = a^2 + \left(\frac{b}{p}\right)^2 + c^2$$

- $a=f(R_1 \text{ \& } R_2)$  function of intrinsic resolution of vertex detector
- $b$  due to multiple scattering in detector
- $c$  due to detector alignment and stability



# Impact parameter

- $\sigma_b = f(\text{vertex layers, distance from main vertex, spatial resolution of each detector, material before precision measurement, alignment, stability})$
- Requirements for best measurement
  - Close as possible to interaction point
  - Maximum lever arm  $R_2 - R_1$
  - Maximum number of space points
  - High spatial resolution
  - Smallest amount of material between interaction point and 1<sup>st</sup> layer
  - Good stability and alignment – continuously measured and correct for 100% detection efficiency
  - Fast readout to reduce pile up in high flux environments

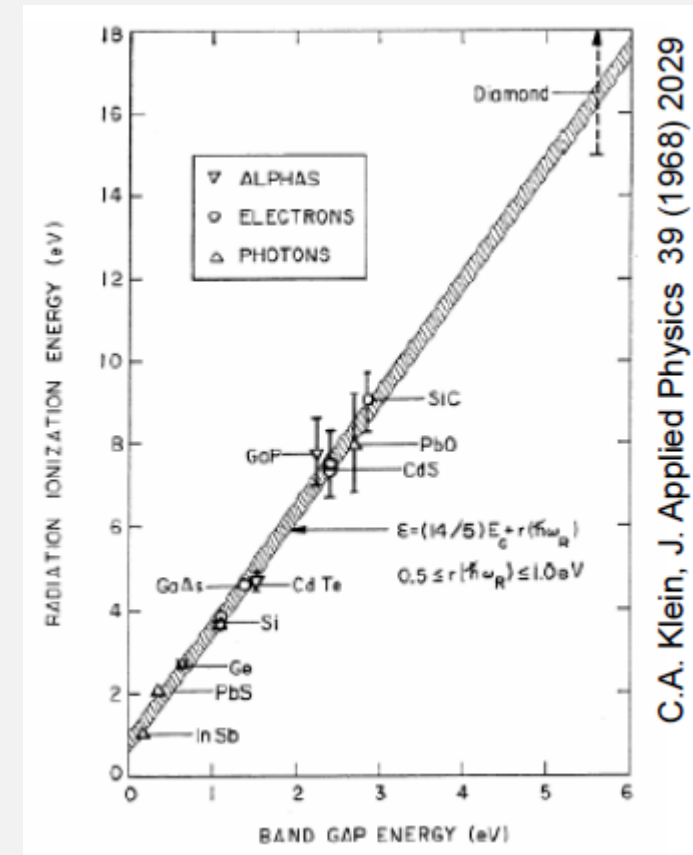
# Basic considerations

Silicon Pixel detectors



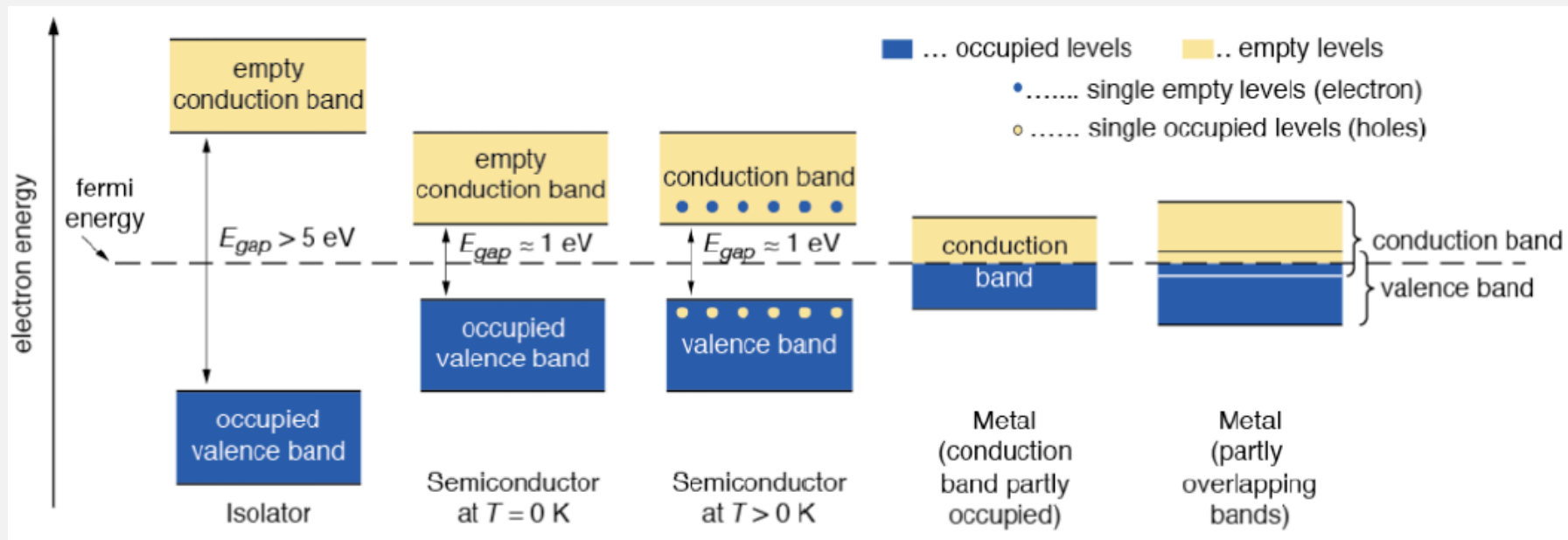
# Why Silicon?

- Semiconductor with moderate bandgap (1.12 eV)
- Energy to create electron/hole pair (signal quanta) = 3.6 eV
  - (c.f. Argon gas = 15 eV)
  - high carrier yield
  - better energy resolution and high signal
- High density and atomic number
  - higher specific energy loss
    - thinner detectors
    - reduced range of secondary particles
    - better spatial resolution
- High carrier mobility → Fast!
  - less than 30 ns to collect entire signal
- **Large experience in industry** with micro-chip technology
- **High intrinsic radiation hardness**



# Energy bands: isolator – semiconductor – metal

In an isolated atom the electrons have only discrete energy levels. In solid state material the atomic levels merge to energy bands. In **metals** the conduction and the valence band **overlap**, whereas in isolators and semiconductors these levels are **separated** by an energy gap (**band gap**). In isolators this gap is large.



# Comparison of different semiconductor materials

Material	Si	Ge	GaAs	GaP	CdTe	Diamond*
Atomic number $Z$	14	32	31+33	31+15	48+52	6
Mass Number $A$ (amu)	28.086	72.61	69.72+74.92	69.72+30.97	112.4+127.6	12.011
Lattice constant $a$ (Å)	5.431	5.646	5.653	5.451	6.482	3.567
Density $\rho$ (g/cm <sup>3</sup> )	2.328	5.326	5.32	4.13	5.86	3.52
$E_g$ (eV) bei 300 K	1.11	0.66	1.42	2.26	1.44	5.47–5.6
$E_g$ (eV) bei 0 K	1.17	0.74	1.52	2.34	1.56	$\approx 6$
rel. permittivity $\epsilon_r = \epsilon/\epsilon_0$	11.9	16.0	12.8	11.1	10.9	5.7
Melting point (° C)	1415	938	1237	1477	1040	3527
eff. e <sup>-</sup> -mass ( $m_n/m_e$ )	0.98, 0.19	1.64, 0.08	0.067	0.82	0.11	0.2
eff. hole mass <sup>+</sup> ( $m_h/m_e$ )	0.16	0.044	0.082	0.14	0.35	0.25

\*usually considered an isolator

Source: <http://www.ioffe.rssi.ru/SVA/NSM/Semicond/> ; S.M.Sze, *Physics of Semicon. Devices* , J. Wiley & Sons, 1981,  
 J. Singh, *Electronic & Optoelectronic Properties of Semiconductor Structures*, Cambridge University Press, 2003  
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# Comparison of different semiconductor materials

Material	Si	Ge	GaAs	GaP	CdTe	Diamond*
eff. density of states in conduction band $n_{CB}$ (cm <sup>-3</sup> )	$3 \cdot 10^{19}$	$1 \cdot 10^{19}$	$4.7 \cdot 10^{17}$	$2 \cdot 10^{19}$		$\approx 10^{20}$
eff. Density of states in valence band $n_{VB}$ (cm <sup>-3</sup> )	$1 \cdot 10^{19}$	$6 \cdot 10^{18}$	$7 \cdot 10^{18}$	$2 \cdot 10^{19}$		$\approx 10^{19}$
Electron mobility $\mu_e$ bei 300 K (cm <sup>2</sup> /Vs)	~1450	3900	8500	< 300	1050	1800
Hole mobility $\mu_h$ bei 300 K (cm <sup>2</sup> /Vs)	~450	1900	400	< 150	100	1200
instrins. charge carrier density at 300 K (cm <sup>-3</sup> )	$1.45 \cdot 10^{10}$	$2.4 \cdot 10^{13}$	$2 \cdot 10^6$	2		$\approx 10^{-27}$
instrins. resistivity at 300 K ( $\Omega$ cm)	$2.3 \cdot 10^5$	47	$\approx 10^8$		$\approx 10^9$	$\geq 10^{42}$
Breakdown field (V/cm)	$3 \cdot 10^5$	$\approx 10^5$	$4 \cdot 10^5$	$\approx 10^6$		$3 \cdot 10^7$
Mean $E$ to create an e-h <sup>+</sup> pair (eV), 300 K	3.62	2.9	4.2	$\approx 7$	4.43	13.25

\*usually considered an isolator

Source: <http://www.ioffe.rssi.ru/SVA/NSM/Semicond/>; S.M.Sze, *Physics of Semicon. Devices*, J. Wiley & Sons, 1981,  
J. Singh, *Electronic & Optoelectronic Properties of Semiconductor Structures*, Cambridge University Press, 2003  
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# Doping intrinsic silicon

A pn junction consists of n and p doped substrates:

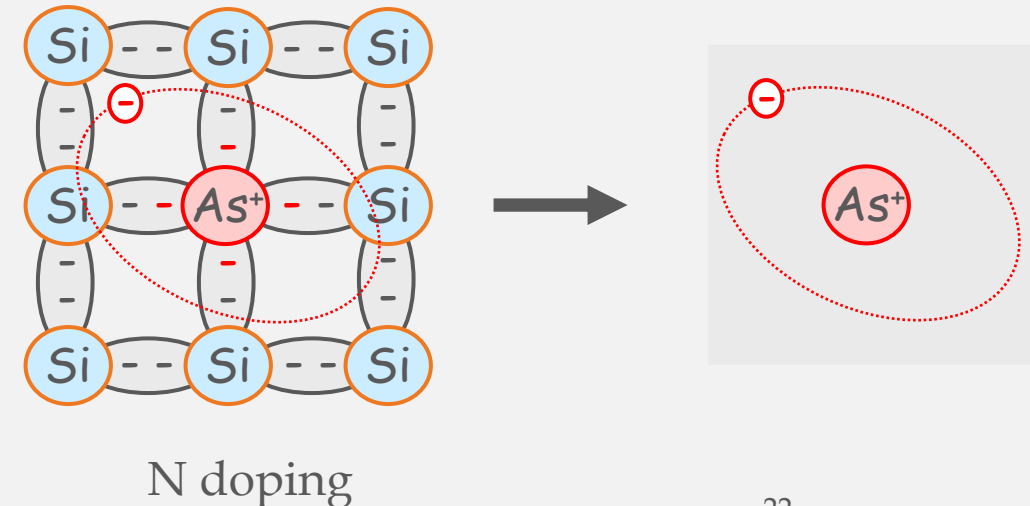
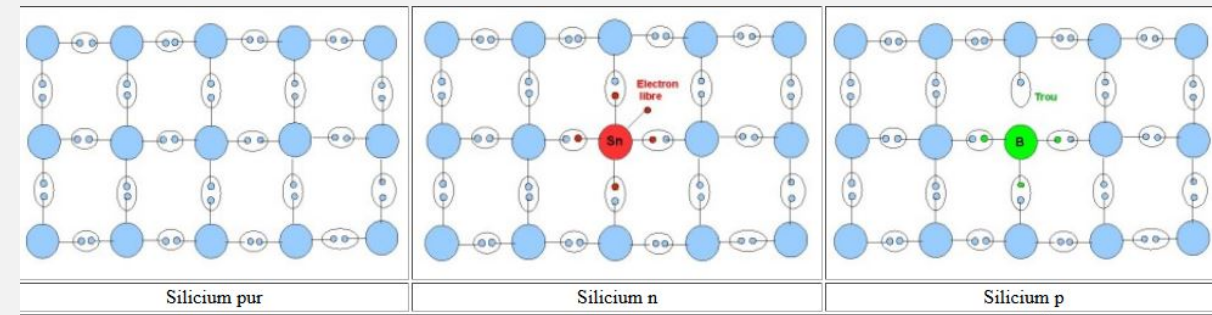
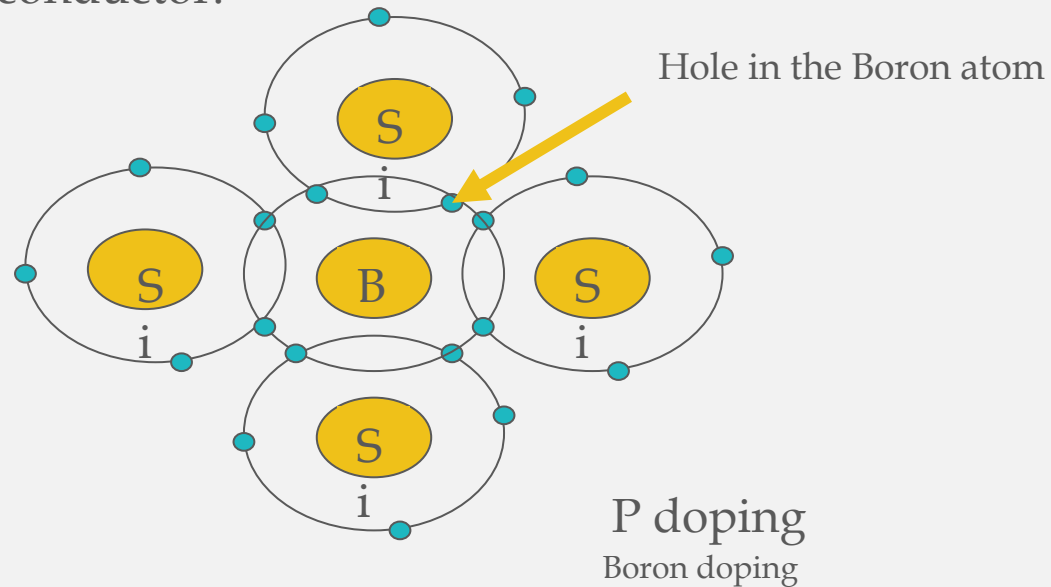
- Doping is the **replacement of a small number of atoms** in the lattice by atoms of **neighboring columns** from the periodic table
- These doping atoms create **energy levels within the band gap** and therefore alter the conductivity.

## Definitions:

- An un-doped semiconductor is called an **intrinsic semiconductor**
  - For each conduction electron exists the corresponding hole.
- A doped semiconductor is called an **extrinsic semiconductor**.
  - Extrinsic semiconductors have a abundance of electrons or holes.

# doping

- In a semi-conductor the electric current is vehiculed by two types of carriers: electrons (negative carriers) and holes (positive carriers).
- N Doping: excess of electron carriers in the semiconductor.
- P doping : excess of holes (deficit of electrons) in the semiconductor.

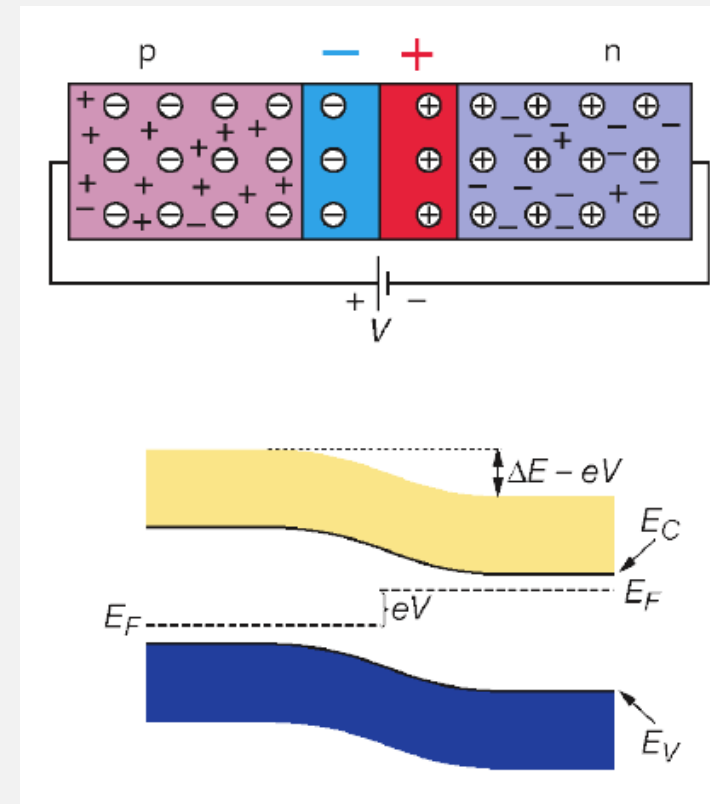


# Operation of a pn-junction with forward bias

Applying an external voltage  $V$  with the anode to p and the cathode to n, e- and holes are refilled to the depletion zone. The **depletion zone becomes narrower** (forward biasing)

## Consequences:

- The potential barrier becomes smaller by  $eV$
- Diffusion across the junction becomes easier
- The current across the junction increases significantly.

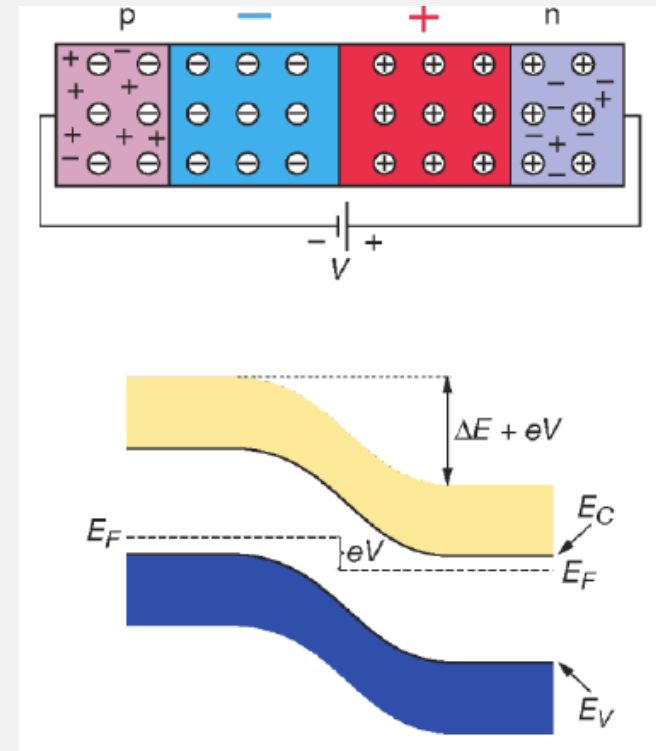


# Operation a pn-junction with reverse bias

Applying an external voltage  $V$  with the cathode to p and the anode to n  $e^-$  and holes are pulled out of the depletion zone.  
The **depletion zone becomes larger** (reverse biasing).

## Consequences:

- The potential barrier becomes higher by  $eV$
- Diffusion across the junction is suppressed.
- The current across the junction is very small (“leakage current”)



➤ This is the way we operate our semiconductor detector!



# Drift velocity and mobility

Drift velocity  
for electrons:

$$\vec{v}_n = -\mu_n \vec{E}$$

and for holes:

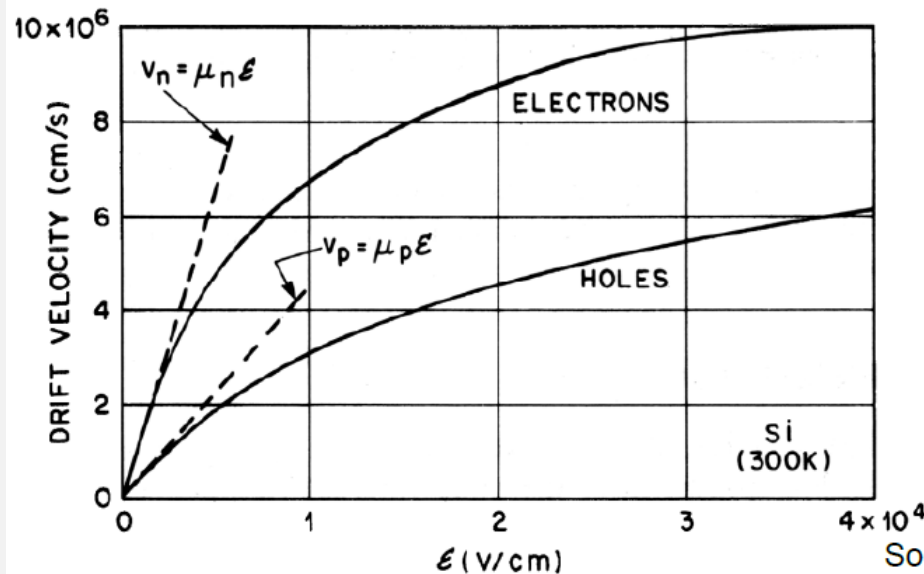
$$\vec{v}_p = -\mu_p \vec{E}$$

Mobility  
for electrons:

$$\mu_n = \frac{e \cdot \tau_n}{m_n}$$

and for holes:

$$\mu_p = \frac{e \cdot \tau_p}{m_p}$$



- $e$  ... electron charge
- $E$  ... external electric field
- $m_n, m_p$  ... effective mass of  $e^-$  and holes
- $\tau_n, \tau_p$  ... mean free time between collisions for  $e^-$  and holes (carrier lifetime)

Source: S.M. Sze, *Semiconductor Devices*, J. Wiley & Sons, 1985

# Width of the depletion zone

Effective doping concentration in typical silicon detector with p<sup>+</sup>-n junction

- $N_a = 10^{15} \text{ cm}^{-3}$  in p<sup>+</sup> region
- $N_d = 10^{12} \text{ cm}^{-3}$  in n bulk

**without external voltage:**

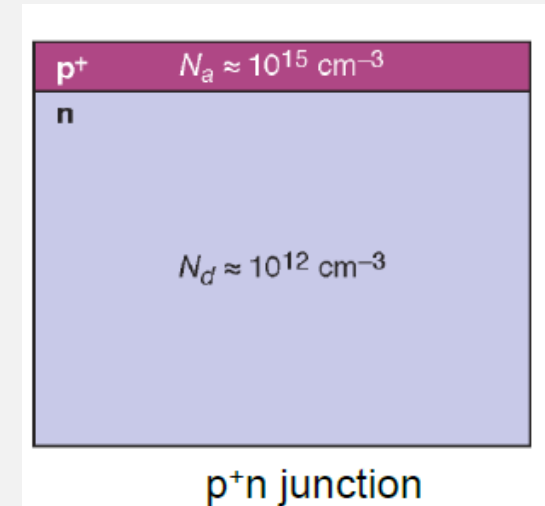
$$W_p = 0.02 \text{ } \mu\text{m}$$

$$W_n = 23 \text{ } \mu\text{m}$$

**Applying a reverse bias voltage of 100 V:**

$$W_p = 0.4 \text{ } \mu\text{m}$$

$$W_n = 363 \text{ } \mu\text{m}$$



Width of depletion zone in n bulk:

$$W \approx \sqrt{2\epsilon_0\epsilon_r\mu\rho|V|}$$

with  $\rho = \frac{1}{e\mu N_{eff}}$

$V$  ... External voltage  
 $\rho$  ... specific resistivity  
 $\mu$  ... mobility of majority charge carriers  
 $N_{eff}$  ... effective doping concentration

# Resistivity

Specific resistivity is a measure of silicon purity

$$\rho = \frac{1}{e(\mu_n n_e + \mu_p n_h)}$$

$n_e, n_h$  ... Charge carrier density for electrons and holes  
 $\mu_n, \mu_p$  ... Mobility for electrons and holes  
 $e$  ... elementary charge

Carrier mobilities:  $\mu_p(\text{Si}, 300\text{K}) \approx 450 \text{ cm}^2/\text{Vs}$   
 $\mu_n(\text{Si}, 300\text{K}) \approx 1450 \text{ cm}^2/\text{Vs}$

The charge carrier concentration in pure silicon (i.e. intrinsic Si) for  $T = 300 \text{ K}$  is:

$$n_e = n_h \approx 1.45 \cdot 10^{10} \text{ cm}^{-3}$$

This yields an intrinsic resistivity of:

$$\rho \approx 230 \text{ k}\Omega\text{cm}$$

# Potential and field inside the junction

- Solving the Poisson equation for an abrupt p-n junction diode

Poisson equation

$$\rho(x) = e \cdot N_{eff} = const$$

$$-\frac{d^2}{dx^2} \phi(x) = \frac{e \cdot N_{eff}}{\epsilon}$$

$$E(w) = 0$$

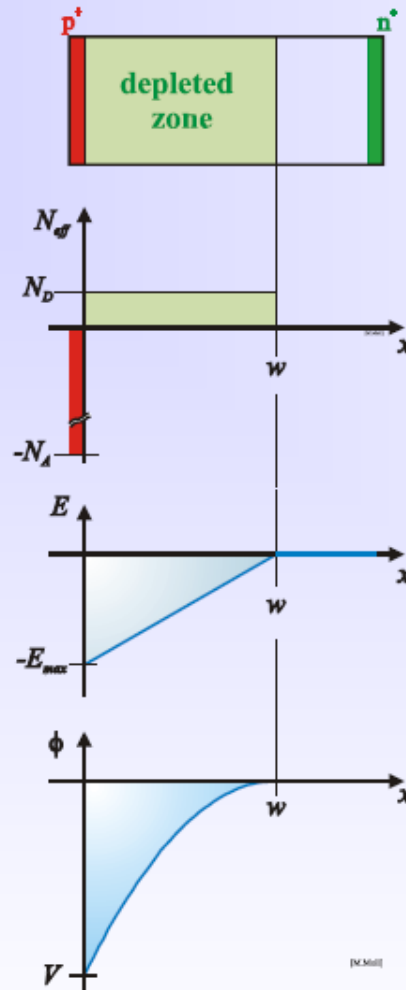
Electric field strength

$$E(x) = -\frac{e \cdot N_{eff}}{\epsilon} \cdot (w - x)$$

$$\phi(w) = 0$$

Electrostatic potential

$$\phi(x) = -\frac{e \cdot N_{eff}}{\epsilon} \cdot \frac{1}{2} \cdot (w - x)^2$$



- electric field strength linear function of depth
- depleted zone growing in depth proportional to  $\sqrt{V}$

$$E_{max} = -\frac{e \cdot N_{eff}}{\epsilon} \cdot w$$

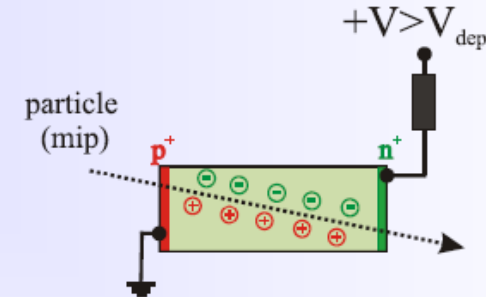
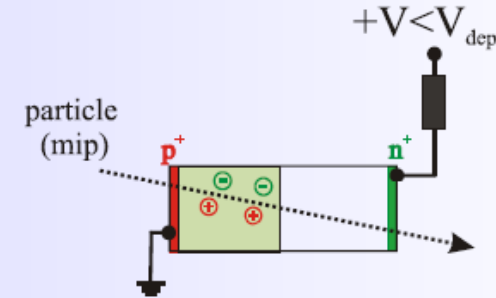
$$w = \sqrt{(-V) \cdot \frac{2 \cdot \epsilon}{e \cdot N_{eff}}}$$

Depletion depth  $w$



# Depletion voltage and sensitive volume

- Below depletion ( $V < V_{dep}$ )
  - Depletion zone  $x_n$  growing with  $w \propto \sqrt{V}$
  - Only charge generated inside depleted volume will be detected
  - Charge generated in 'neutral zone' (field free zone) will recombine
- Depletion Voltage  $V_{dep}$ 
  - Sensor depleted of free charge carriers
  - Electric field throughout complete device
  - Complete sensor volume sensitive (active)
  - Example:
    - $d = 300 \mu\text{m}$
    - $N_{eff} = [P] = 1.5 \times 10^{12} \text{ cm}^{-3}$   
( $\rho \approx 3 \text{ k}\Omega\text{cm}$ )
    - $V_{dep} \approx 100 \text{ V}$
- Full charge collection only for ( $V > V_{dep}$ )



depletion voltage  $V_{dep}$       detector thickness  $d$

$$V_{dep} = \frac{e}{2 \cdot \epsilon} \cdot |N_{eff}| \cdot d^2$$

effective space charge density  $N_{eff}$

# Depletion voltage determination

## □ “ Full Depletion Voltage $V_{FD}$ ”

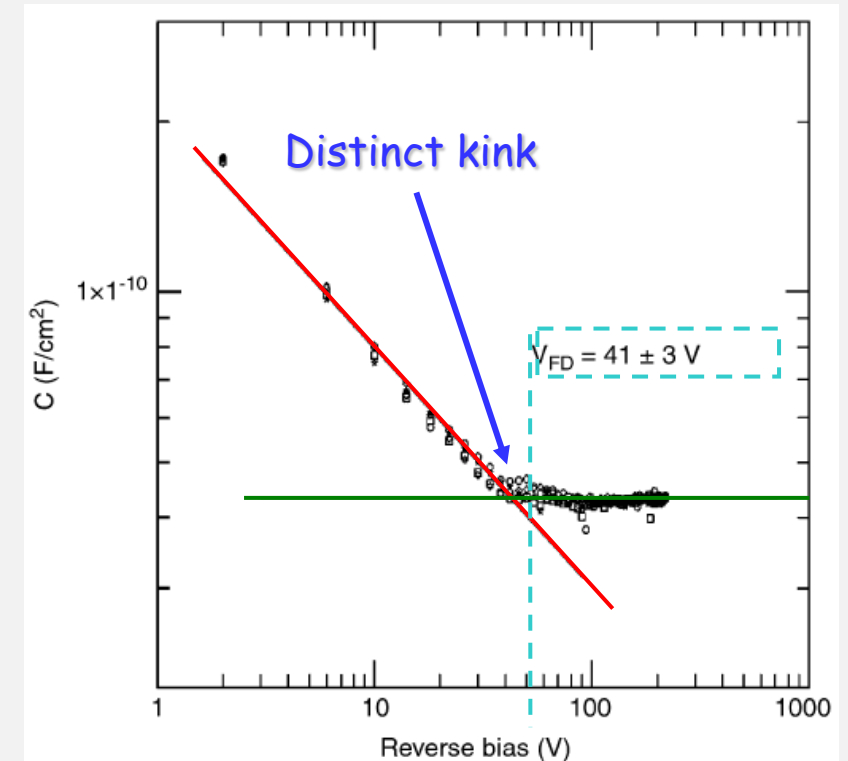
### ▪ Definition

- For a semiconductor detector to operate, it is necessary to create a carrier depletion layer in the pn-junction area.
- The potential where this condition is reached is the “ Full Depletion Voltage ” ( $V_{FD}$ ).

### ▪ Determination of $V_{FD}$ from CV curves

The evaluation of  $V_{FD}$  can be performed by the double linear fit of the logarithm of the capacitance versus the logarithm of the voltage (CV method).

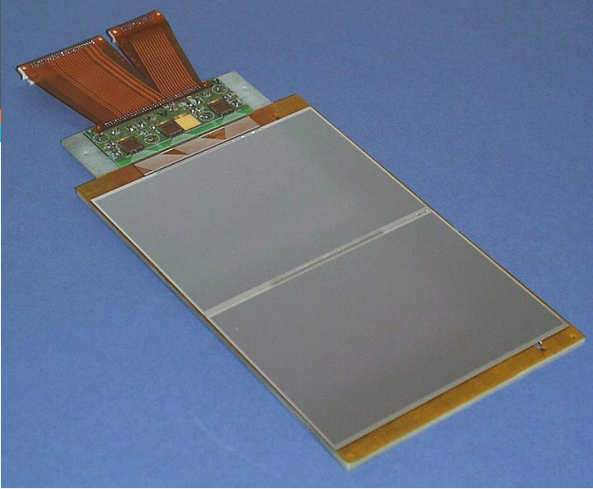
The frequency is generally taken around 1-10 kHz.



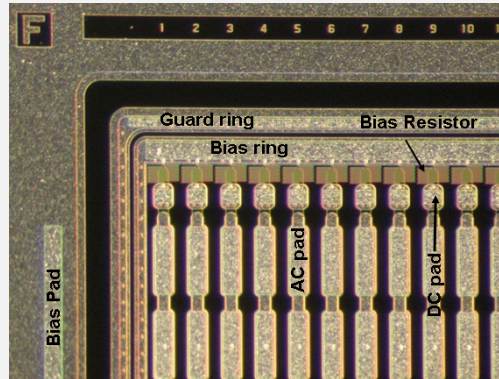
# Design considerations

## " The Key Parameters "

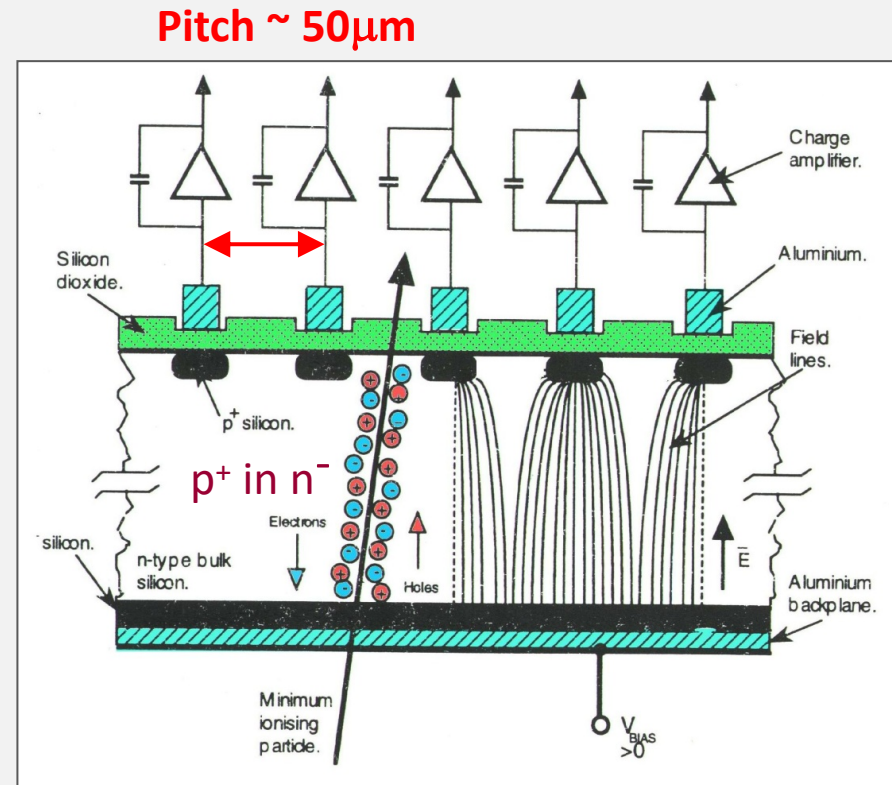
# Micro-strip Silicon Detectors



Highly segmented silicon detectors have been used in Particle Physics experiments for nearly 30 years. They are favourite choice for Tracker and Vertex detectors  
(high resolution, speed, low mass, relatively low cost)



Main application: detect the passage of ionizing radiation with high spatial resolution and good efficiency.  
Segmentation → position



Resolution  $\sim 5\mu\text{m}$



# Silicon pixel sensor principle

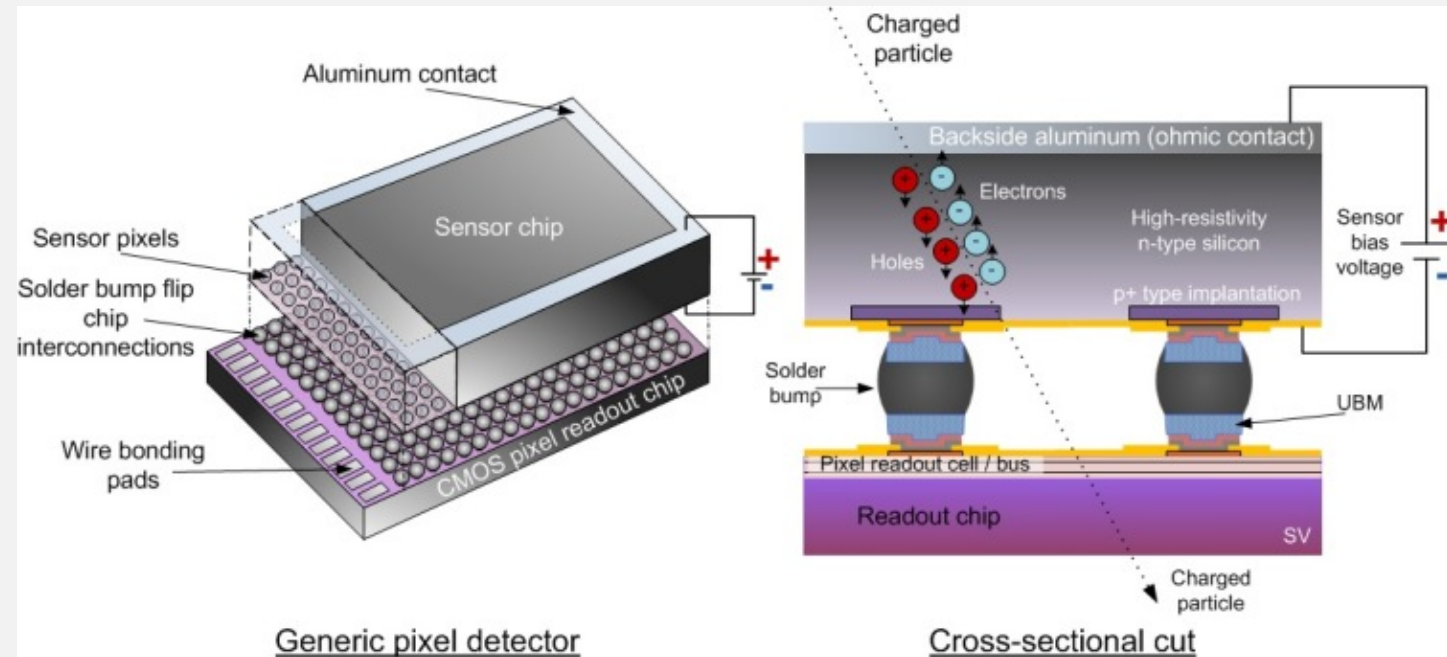
- Operation

- Multi-PN diode in reverse bias
- Depletion layer start from junction
- Particle ionises Si, producing e/h pairs
- e/h pairs drift in E-Field to electrodes

- Properties

- Thickness from 50 to 300  $\mu\text{m}$
- → signal 24000 e/h pairs ( $80 \text{ e}/\mu\text{m}$ )
- High segmentation required

- Advantage of using silicon
  - Low ionisation energy
  - Fast signal collection



- Relevant parameters for performance

- Leakage current
- Depletion voltage ( $V_{\text{dep}}$ ) Operational Voltage
- Heat load

# Noise

- Depends upon detector capacitance and reverse current
- Depends upon electronics design
- Function of signal shaping time
- Lower capacitance  $\Rightarrow$  lower noise
- Faster electronics  $\Rightarrow$  noise contribution from reverse current less significant

# Constructing a detector

Let's make a simple calculation for silicon:

- mean ionization energy  $I_0 = 3.62 \text{ eV}$ .
- mean energy loss per flight path of a mip  $dE/dx = 3.87 \text{ MeV/cm}$

Assuming a detector with a thickness of  $d = 300 \text{ }\mu\text{m}$  and an area of  $A = 1 \text{ cm}^2$

- **Signal of a mip in such a detector:**

$$\frac{dE/dx \cdot d}{I_0} = \frac{3.87 \cdot 10^6 \text{ eV/cm} \cdot 0.03 \text{ cm}}{3.62 \text{ eV}} \approx 3.2 \cdot 10^4 e^- h^+ - \text{pairs}$$

- **Intrinsic charge carrier in the same volume ( $T = 300 \text{ K}$ )**

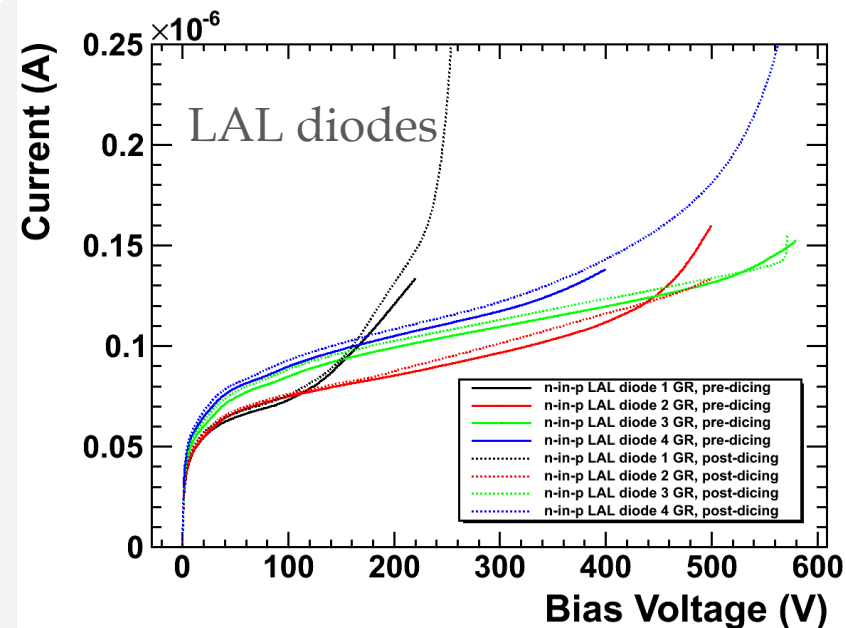
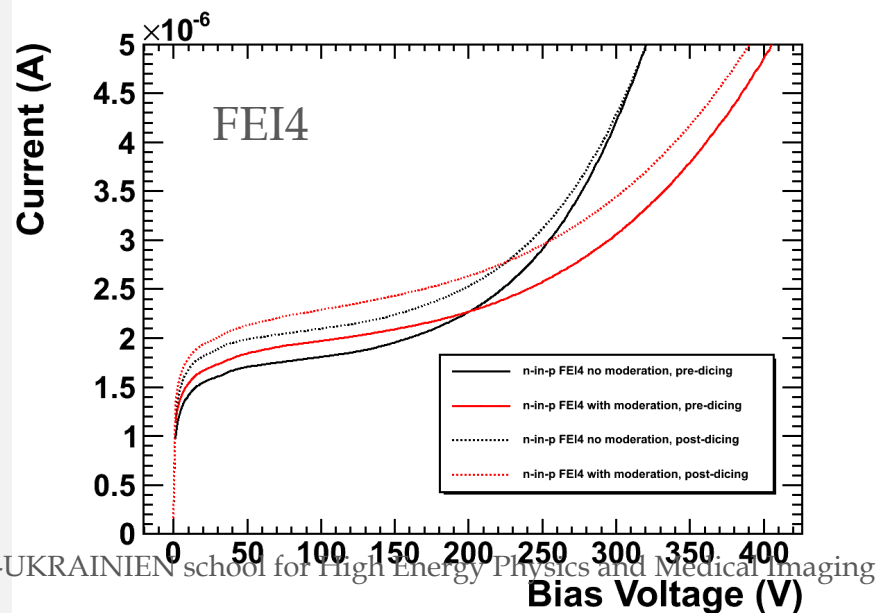
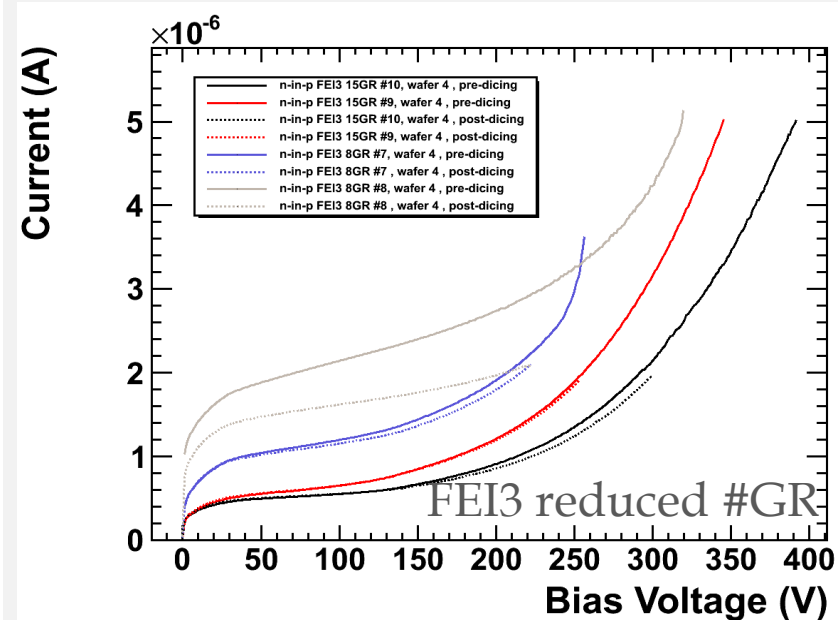
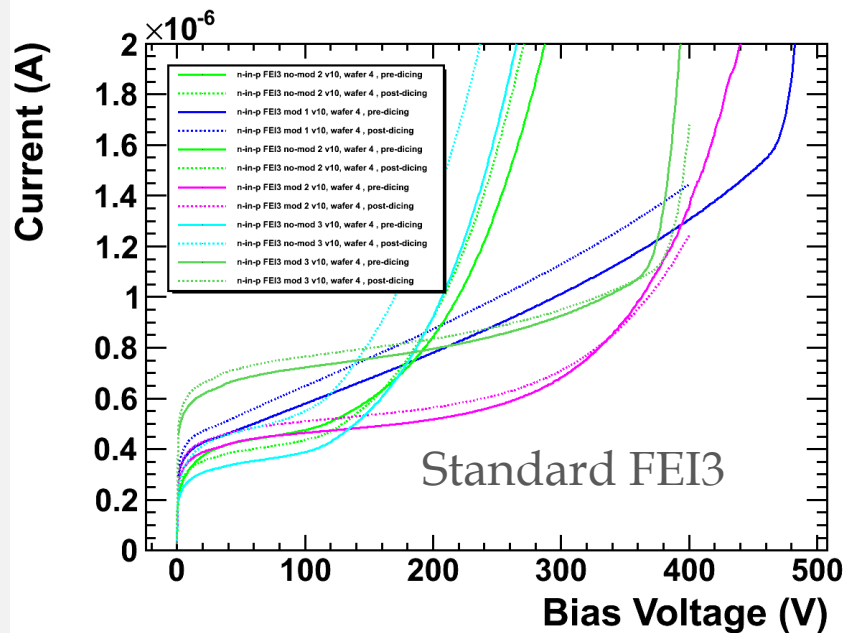
$$n_i \cdot d \cdot A = 1.45 \cdot 10^{10} \text{ cm}^{-3} \cdot 0.03 \text{ cm} \cdot 1 \text{ cm}^2 \approx 4.35 \cdot 10^8 e^- h^+ - \text{pairs}$$

**Result: the number of thermal created e-h+-pairs (noise is four orders of magnitude larger than the signal).**

We have to remove the charge carriers

→ depletion zone in inverse biased **pn junctions**

# Reverse current : characteristics on p-type wafer: a selection criteria



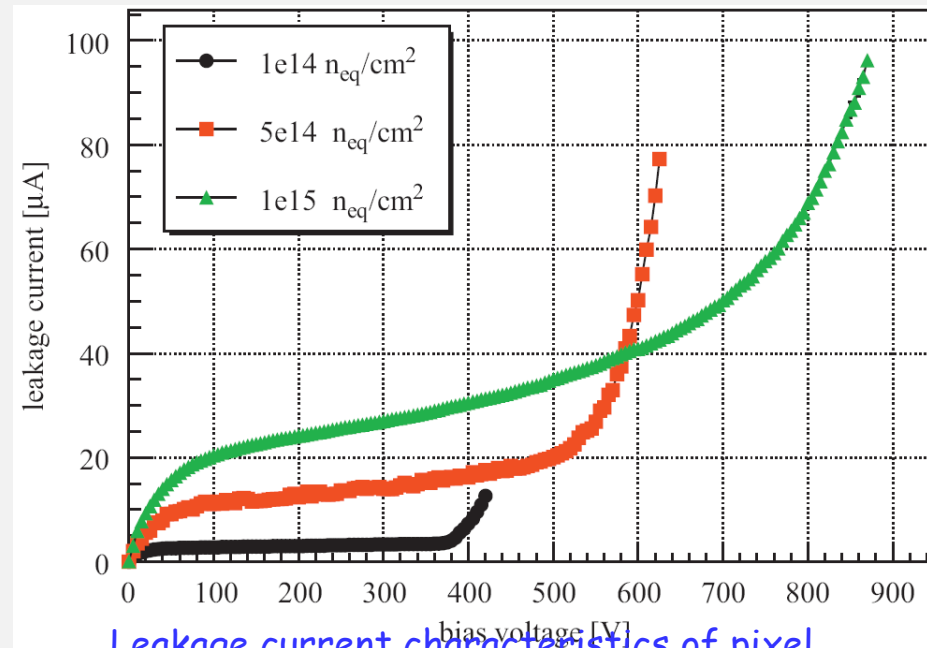
# Leakage current and breakdown of irradiated devices

## □ “ Break-Down Voltage $V_{BD}$ ”

- The breakdown voltage is defined as the highest measured voltage for a defined normalized leakage current.

## □ “ Leakage Current ”

- The leakage current is a sensitive and effective signature for quality assurance and control. It originates from the introduction of generation/recombination centers.



Leakage current characteristics of pixel sensors for different proton fluences.



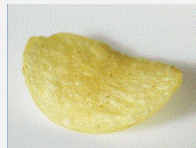
# Design considerations

## Thin planar process

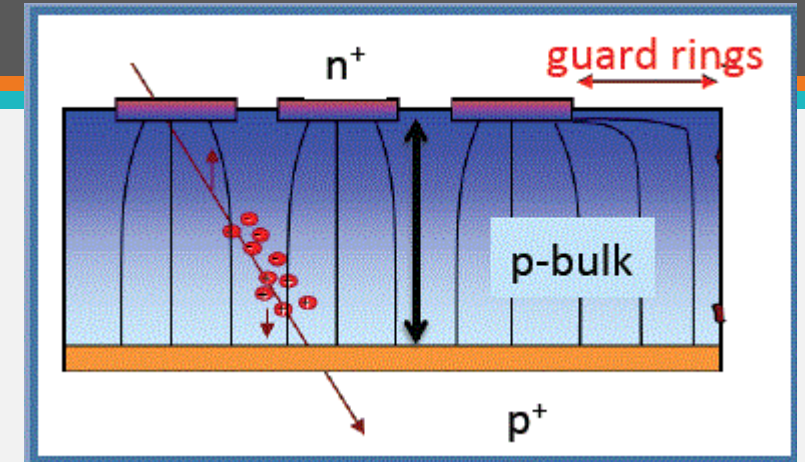
- ❖ Low leakage current (before and after irradiation)
- ❖ Low operational voltage
- ❖ Less power consumption
- ❖ Short drift path → less trapping
- ❖ Higher electric field at low  $V_{bias}$  (fast collection time)
- ❖ Less material (good for tracking)

## Drawback

- ❖ Smaller initial signal
- ❖ Thinning cost extra money
- ❖ Handle with care

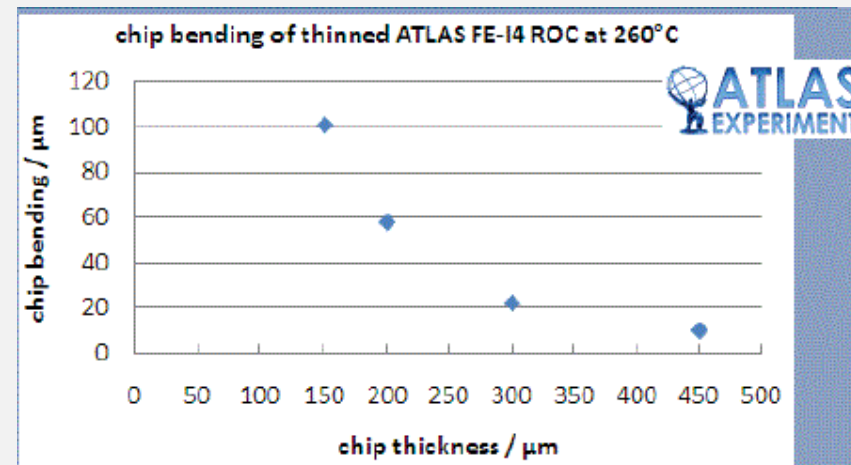


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## Readout chip connection to **thin** sensor

- Can lead to internal stress
- Can introduce disconnected bumps (open)
- Alignment inaccuracies



# How much charge can be collected : The Charge signal

## ■ Collected Charge for a Minimum Ionizing Particle (MIP)

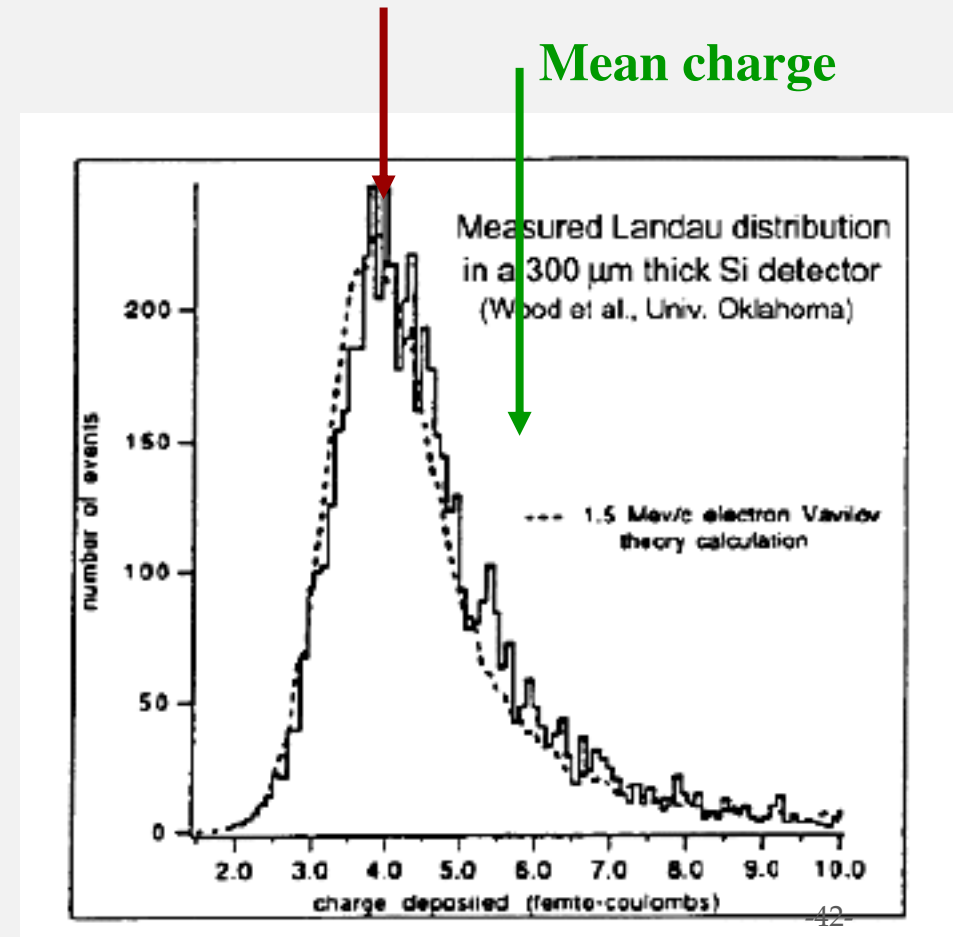
- Mean energy loss  
 $dE/dx$  (Si) = 3.88 MeV/cm  
 $\Rightarrow$  116 keV for 300  $\mu$ m thickness
- Most probable energy loss  
 $\approx 0.7 \times \text{mean}$   
 $\Rightarrow$  81 keV
- 3.6 eV to create an e-h pair  
 $\Rightarrow$  108 e-h /  $\mu$ m (mean)  
 $\Rightarrow$  ~80 e-h /  $\mu$ m (most probable)
- Most probable charge (300  $\mu$ m)

$\approx 24000$  e       $\approx 3.6$  fC

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Medical Imaging

Most probable charge  $\approx 0.7 \times$  mean

Mean charge



## A few considerations on charge collection in case of trapping (Irradiation)

- Charge collection in pixel detectors:

The total charge collected by a detector , Q, can be written as :  $Q = Q_0 \cdot \theta \cdot \eta$

$\theta$  is the geometrical factor and  $\eta$  the trapping factor

For a planar pad detector with no segmentation, the weighting field  $E_w$  is  $1/d$  (ramo\*\*) and the induced current by a moving charge  $Q'_0(t)$  is :

$$i(t) = Q'_0(t) \cdot E_w \cdot v_{dr}(t) = Q'_0(t) \cdot \frac{1}{d} \cdot v_{dr}(t)$$

$v_{dr}(t)$  is the carrier drift velocity that depends on the detector internal field, and  $Q'_0(t)$  is affected by trapping:  $Q'_0(t) = Q_0 \cdot e^{-\frac{t}{\tau_t}}$

The total collected charge can be written :  $Q(t) = \int_0^{t_{dr}} Q_0 \cdot e^{-\frac{t}{\tau_t}} \cdot \frac{1}{d} \cdot v_{dr}(t) \cdot dt$

$\tau_t$  is the trapping time constant that is inversely proportional to the trap concentration.  $t_{dr}$  is the drift time or carrier transient time

It could be rewritten as :

For constant electric field, we assume a constant drift velocity  $v_{dr}(t) = V_{dr}$ , we have :

$w$  is the full detector depletion depth and  $w = V_{dr} \cdot t_{dr}$  is used for planar detectors

$$Q = Q_0 \cdot \underbrace{\frac{w}{d}}_{\theta} \cdot \frac{\tau_t}{t_{dr}} \cdot \underbrace{(1 - e^{-t_{dr}/\tau_t})}_{\eta}$$

# Charge collection distance for large trapping

- If  $Q_0$  is the charge deposited by a MIP particle in a detector of thickness  $d$  (**in microns**), then  $Q_0$  can be expressed in terms of number of electrons:  $Q_0(t) = 80.d$

- So,  $Q(t)$  in electrons could be written as : 
$$Q(t) = 80.w.\frac{\tau_t}{t_{dr}} \left( 1 - e^{-\frac{t_{dr}}{\tau_t}} \right)$$

- For situation with large trapping , such as HL-LHC, up to  $10^{16} \text{ n}_{eq}/\text{cm}^2$  , we have :  $\frac{\tau_t}{t_{dr}} \ll 1$

Thus : 
$$Q \cong 80.w.\frac{\tau_t}{t_{dr}} = 80.v_{dr}.\tau_t \equiv 80.d_{CCE}$$

- $d_{CCE}$  is expressed in microns and is defined as a **charge collection distance or trapping distance**
- Note that  $d_{CCE}$  is equal to **20 microns** for  $10^{16} \text{ n}_{eq}/\text{cm}^2$  if  $\tau_t$  is 0.2 ns,  $v_{dr}=10^7 \text{ cm/s}$
- Conclusion : it doesn't matter if a detector with thickness  $d > 50 \mu\text{m}$  is fully depleted or not; the collected charge is about the same if  $d \geq 50 \mu\text{m}$*

It is in the order of **1600 electrons** for  $10^{16} \text{ n}_{eq}/\text{cm}^2$  .

## Low or moderate fluences

- For low or moderate fluences, where:  $\frac{\tau_t}{t_{dr}} \gg 1$

$$Q = 80w \cdot \left( 1 - \frac{1}{2} \cdot \frac{t_{dr}}{\tau_t} \right)$$

- In this case, the detector depletion depth is still an important factor in the collected charge, which depends on the voltage :

$$w = \sqrt{\frac{2\varepsilon \cdot \varepsilon_0 \cdot V}{e \cdot N_{eff}}}$$

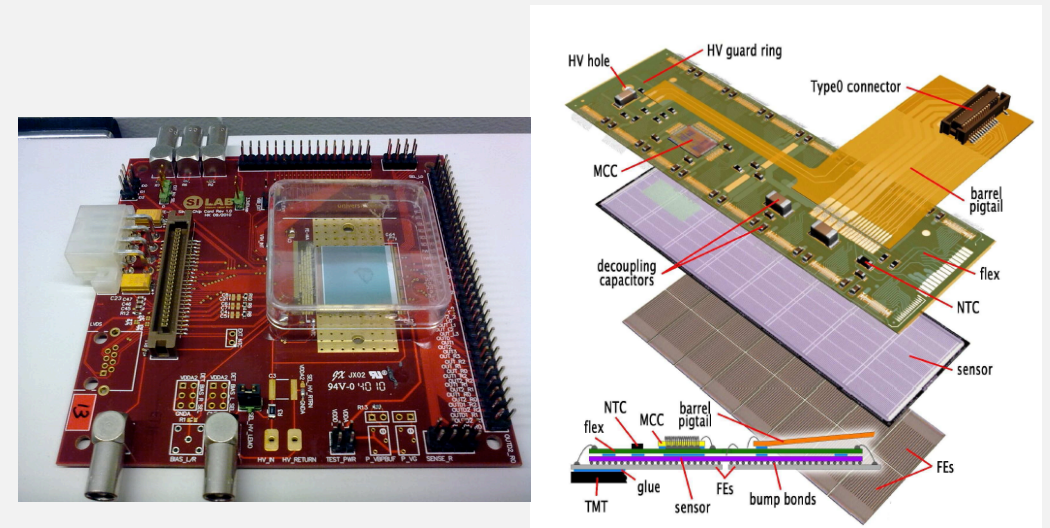
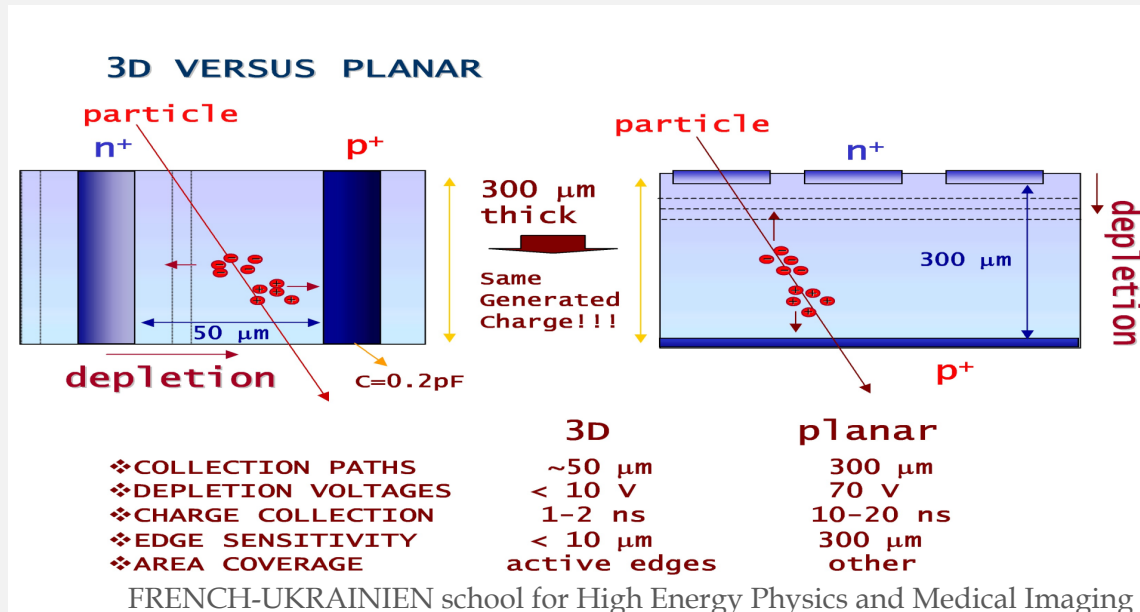
Good Idea

- The idea is to decouple the drift length and the detector thickness ==➔ Go toward 3D detectors where condition is satisfied (column separation 40 to 400  $\mu\text{m}$ ) and the total induced charge by mip is still 80 d ➔ 24000 electrons for 300  $\mu\text{m}$ .



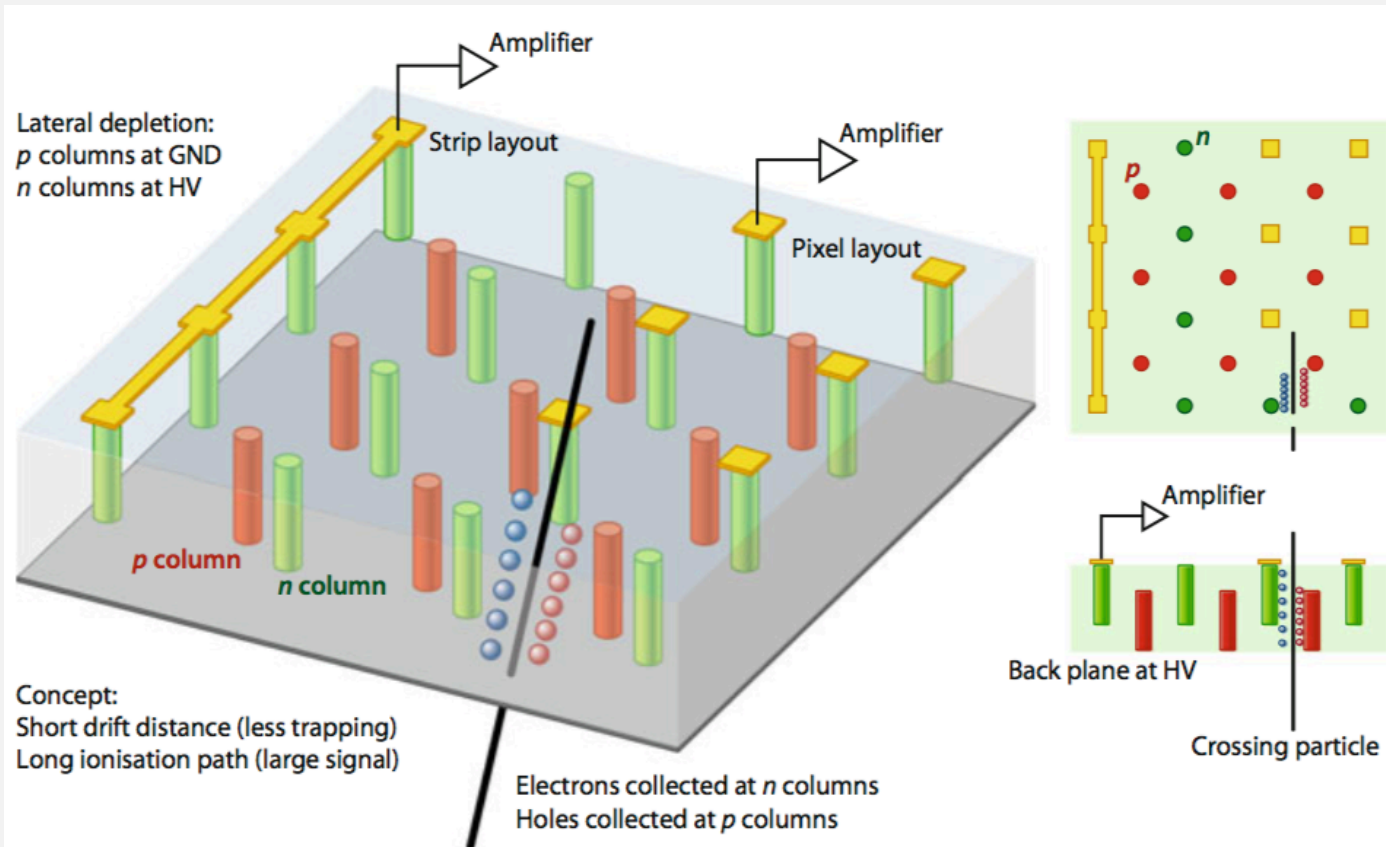
# Improving charge collection efficiency

- The weighting field in 3D detector is given by :  $Q_{3D} = 80 \cdot \frac{d}{\lambda_{cc}} d_{CCE}$  (very high fluences)
- $\lambda_{cc}$  is the electrode spacing
- It is clear that the collected charge in a 3D detector is amplified by  $d/\lambda_{cc}$ , as compared to planar sensors.
- For a 300  $\mu\text{m}$  thick sensor with 50  $\mu\text{m}$  column spacing, the amplification is 6 !*

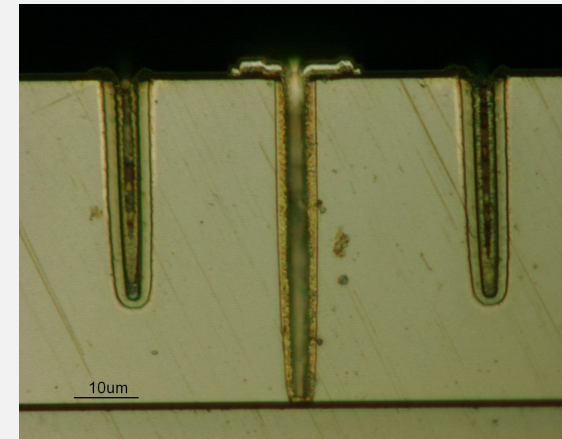
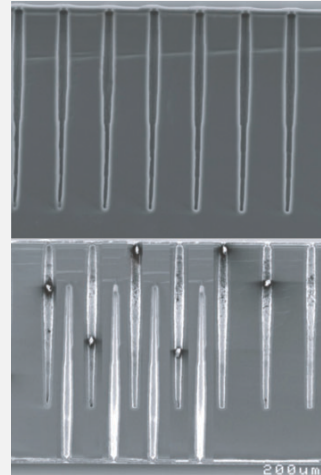


ATLAS FEI4 module

# 3D Principle



- Double or single type
- small pixel cells  $25 \times 100 \mu\text{m}^2$ 
  - We need narrow columns!
- High column depth/width is better...



# Silicon pixel detector for ATLAS

challenges and Constraints



# Planar Technology

The **planar process** is a manufacturing process used in the semiconductor industry to build individual components of a transistor, and in turn, connect those transistors together.

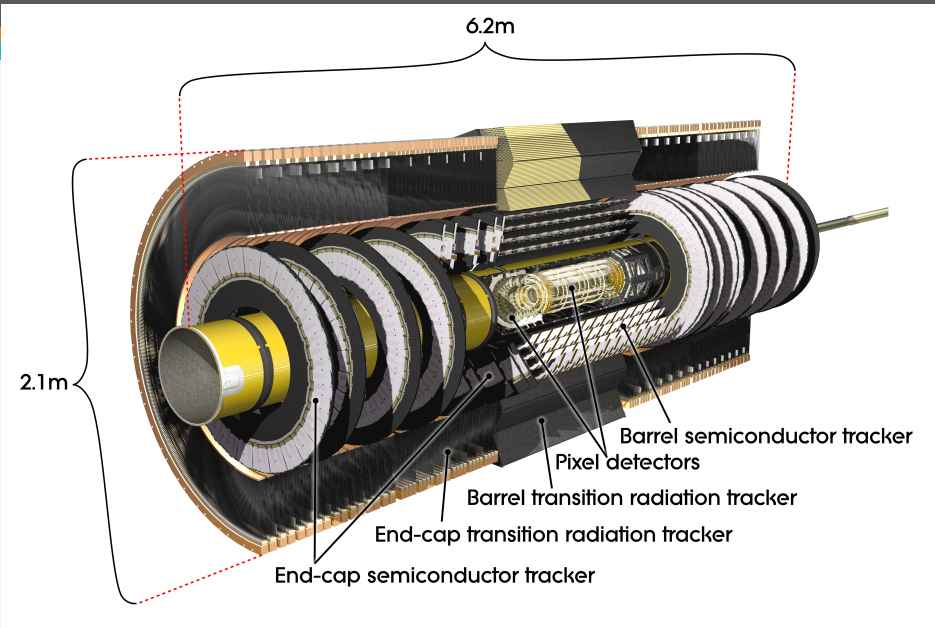
It is the primary process by which modern integrated circuits are built. **The process was developed by Jean Hoerni, while working at Fairchild Semiconductor, with a first patent issued 1959.**

The key concept was to view a circuit in its two-dimensional projection (a plane), thus allowing the use of **photographic processes concepts** such as film negatives to mask the projection of light exposed chemicals. This allowed the use of a series of exposures on a substrate (Silicon) to create **silicon oxide (insulators) or doped regions (conductors)**. Together with the use of **metallization** (to join together the integrated circuits), and **the concept of p-n junction isolation (from Kurt Lehovec)**, the researchers at Fairchild were able to create circuits on a single silicon crystal slice (a wafer) from a mono-crystalline silicon bar.

The process involves the basic procedures of ( $\text{SiO}_2$ ) oxidation,  $\text{SiO}_2$  etching and heat diffusion.

The final steps involves oxidizing the entire wafer silicon dioxide with an  $\text{SiO}_2$  layer, etching contact vias to the transistors, and depositing a covering metal layer over the oxide, thus connecting the transistors without manually wiring them together.

# Challenges for a new inner tracker



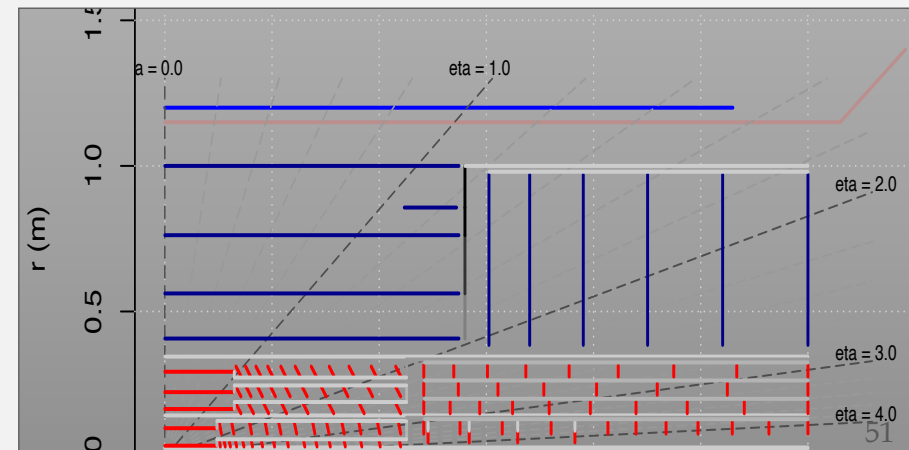
High Luminosity LHC conditions will affect  
Drastically the current design

- Peak Luminosity will increase by **a factor 5-7** :  $5-7 \cdot 10^{34}$
- Average Pileup : **a factor of 8**  $\langle \mu \rangle \sim 200$
- Integrated luminosity : **a factor 10** :  $3000 \text{ fb}^{-1}$
- Radiation hardness : **a factor 20** :  $2 \times 10^{16} \text{ n}_{\text{eq}}/\text{cm}^2$

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## Upgrade of ATLAS phase 0

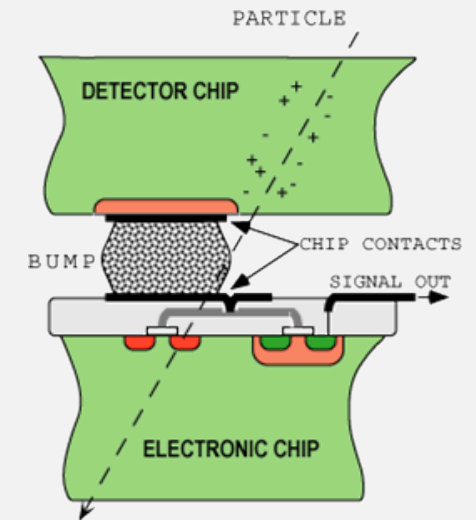
- Same or better performance required !
  - Use quite “transparent tracker” Low X0
  - Provide V. High granularity sensors
  - Go to deep submicron technologies -65 nm CMOS
  - Raise the transfer data band-width (5Gps)
  - Improve the layout and mechanics



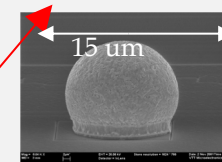
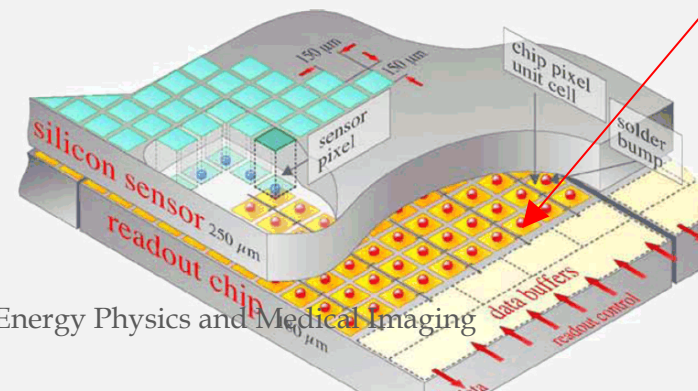


# Hybrid Pixel Detectors used in ATLAS for 900 MChannels

- Hybrid Active Pixel Sensors
  - segment silicon to diode matrix with high granularity (50x50  $\mu\text{m}^2$  or 50x125  $\mu\text{m}^2$ )  
( $\Rightarrow$  true 2D, no reconstruction ambiguity)
  - readout electronic with same geometry (every cell connected to its own processing electronics)
  - connection by “bump bonding”
  - requires sophisticated readout architecture
  - Hybrid pixel detectors will be used in LHC experiments: ATLAS, ALICE, CMS and LHCb



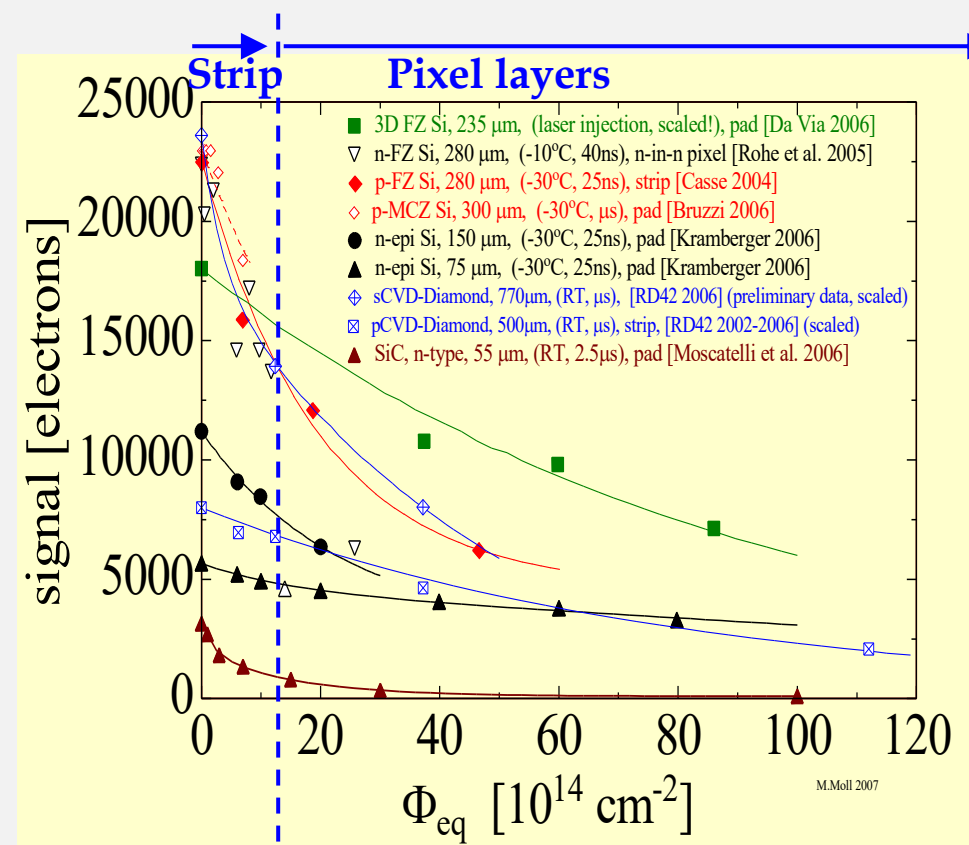
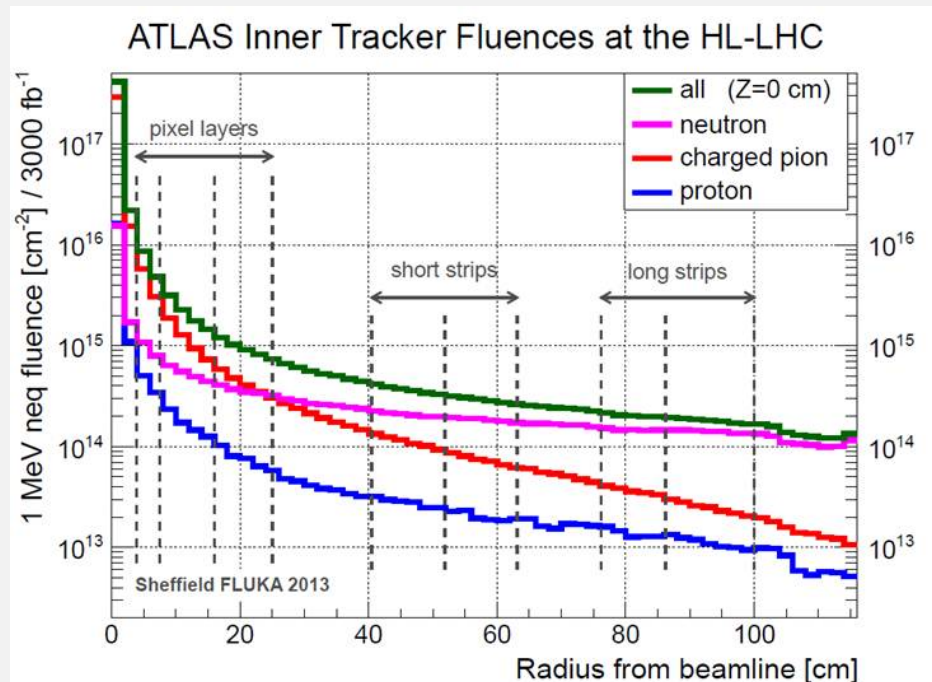
Flip-chip technique



Solder Bump: Pb-Sn

# Fluences and signal collection

(HL-LHC conditions )



# Effects of Radiation in pixel sensors

## 1. Radiation induced leakage current

independent of impurities; every 7°C of temperature reduction halves current

⇔ cool sensors to  $\approx -25^\circ\text{C}$

## 2. “type inversion” from n to p-bulk

⇔ increased depletion voltage

oxygenated silicon helps (for protons);

n<sup>+</sup>-in-n-bulk or n<sup>+</sup>-in-p-bulk helps

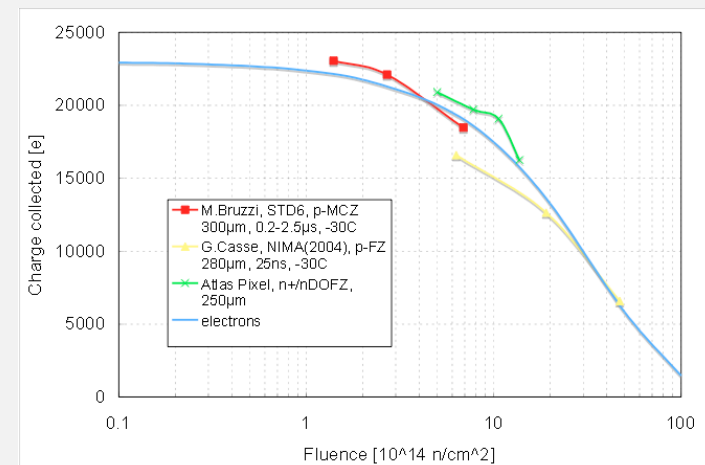
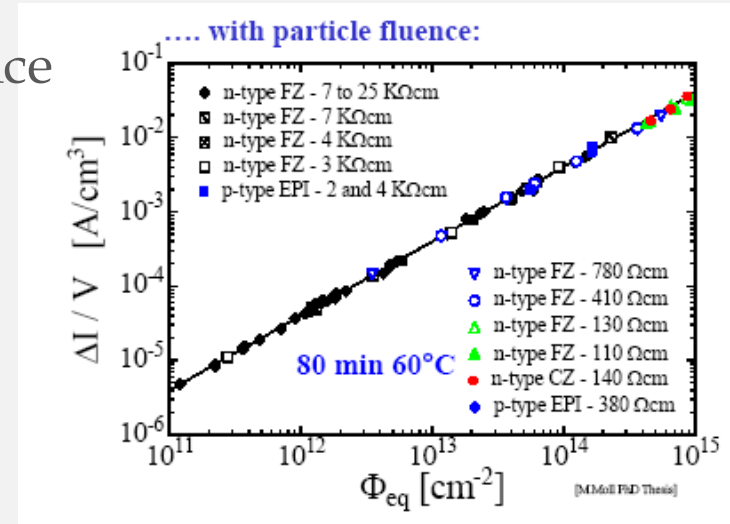
## 3. Charge trapping

the most dangerous effect at high fluences

⇔ collect electrons rather than holes

⇔ reduce drift distances

Current/fluence

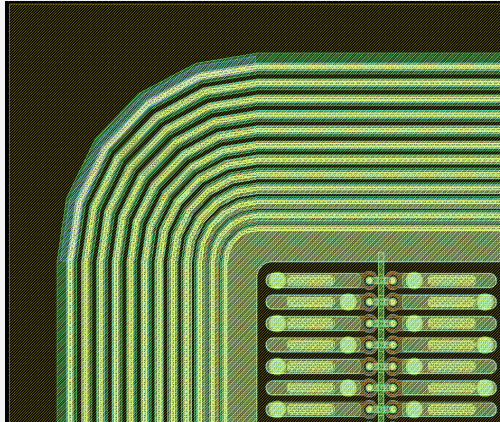


# Overview of R&D activity

- How to improve the actual design ?
- Telescope Tool
- TCAD simulation Tool
- SIMS Tool



# Summary of R&D contributions within ATLAS



## Goals

- Performance : evaluate & improve sensor design for radiation tolerance up to  $3 \times 10^{15} n_{eq}/cm^2$  fluence.
- Productions : work on various wafer productions (Cis, VTT Advacam & FBK)

### Active Edge

- Smaller Edge size
- Higher Efficiency at Edge region

### Biasing Structure

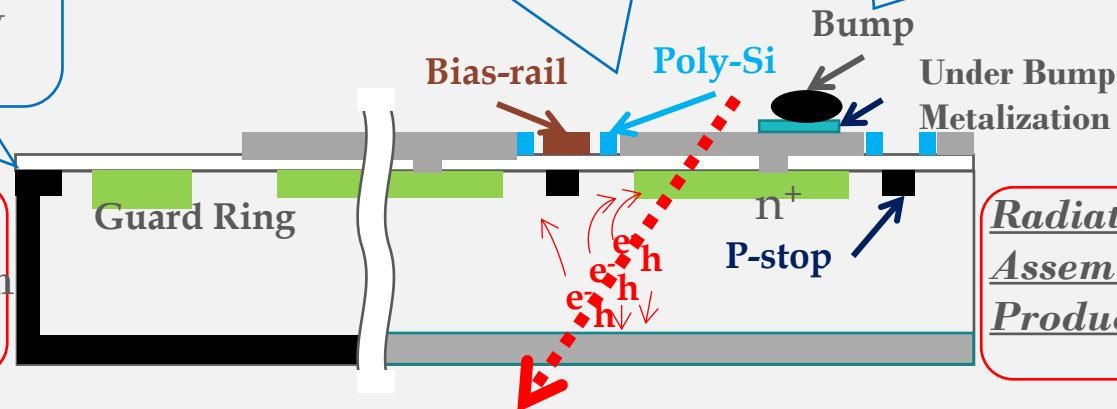
- Optimize Biasing structure
- Higher Efficiency at Pixel boundary

### Bumpbonding

- Better UBM material
- Higher Flip-Chip Yield

### Simulation

- Doping concentration
- Charge collection



Radiation Tolerance  
Assembly  
Production...



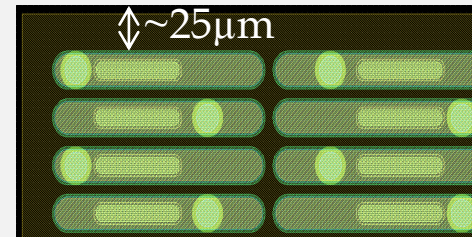
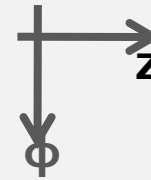
# Typical examples of VTT-LAL Edge design



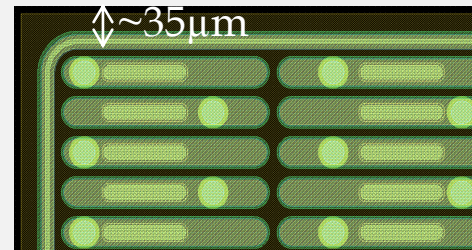
- VTT OMEGAPIX (SlimEdge & Edgeless) designs:

1. 12 GR + BR
2. 1GR + BR
3. 0 GR + BR
4. 0 GR & 0 BR
5. 1GR & no BR

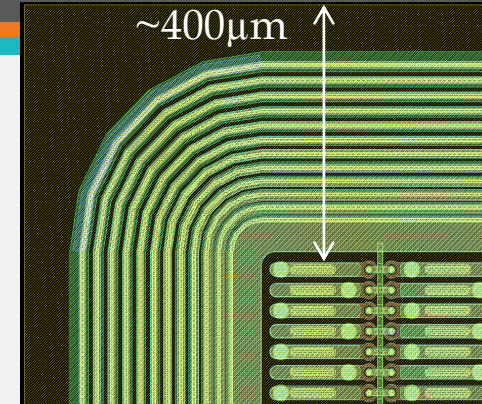
T. Rashid, PhD thesis  
2019



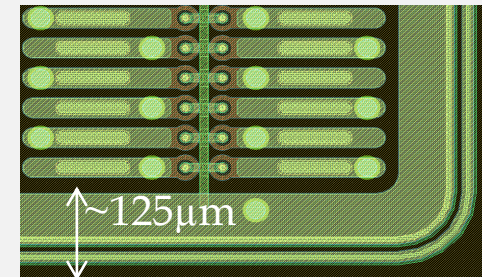
0GR + 0BiasRing



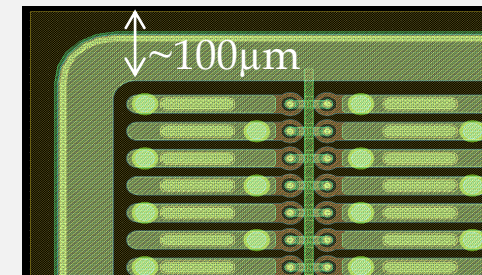
0BiasRing+1GR



12GR+BR



1GR + BiasRing

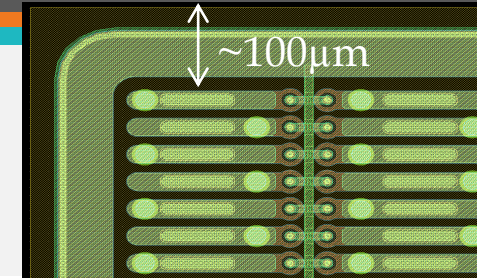
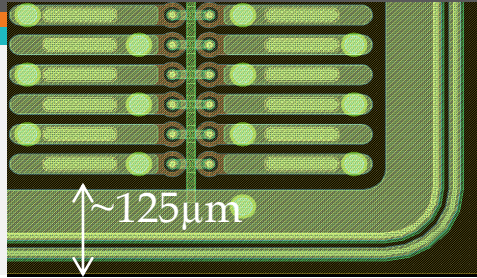


0GR+ BiasRing

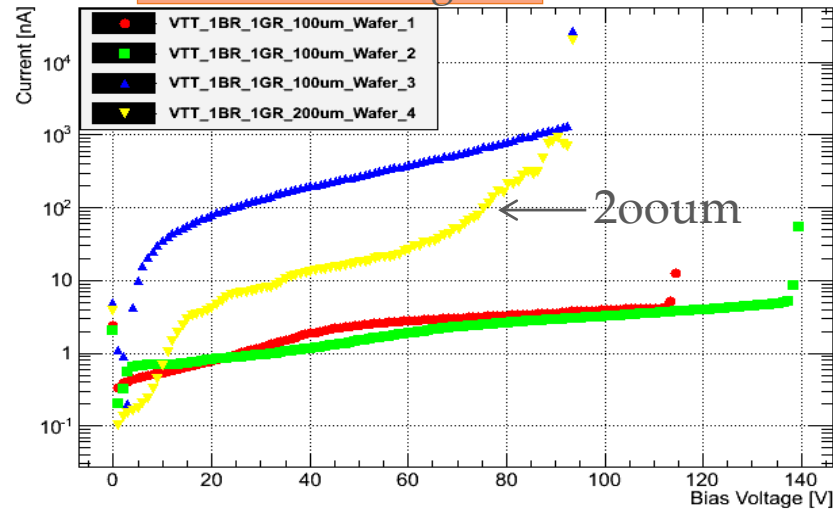
- All designs:

- Active Area:  $4800 \times 3360 \mu\text{m}^2$
- Array:  $96 \times 24 (\phi, z)$
- Pixel Size :  $35 \times 200 \mu\text{m}^2$
- Thickness: 50 to  $200 \mu\text{m}$

# VTT Slim Edge:

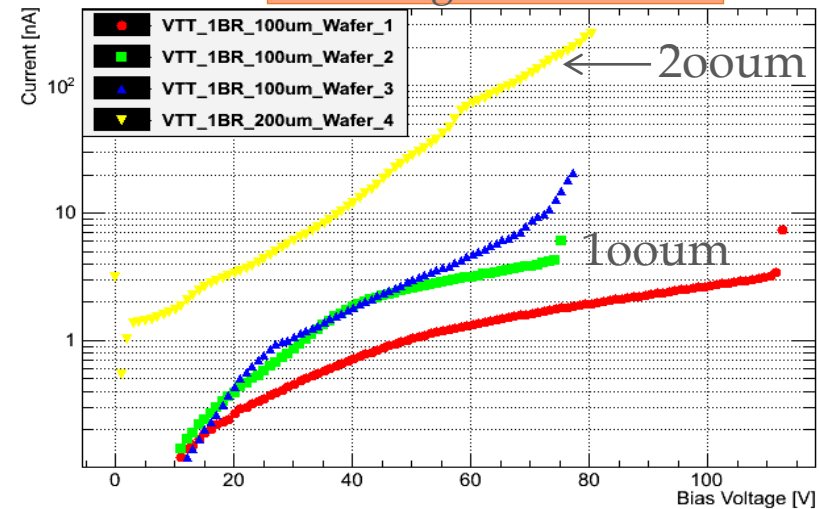


1GR + BiasRing



VTT: BiasRing + 1GR  
100-200μm thickness  
V<sub>bd</sub> ~ 100V to 140V

BiasRing + 0 GR



VTT: BiasRing  
100-200μm thickness  
V<sub>bd</sub> ~ 75V to 120V

# ADVACAM Active Edge Sensors



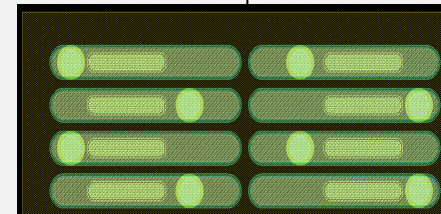
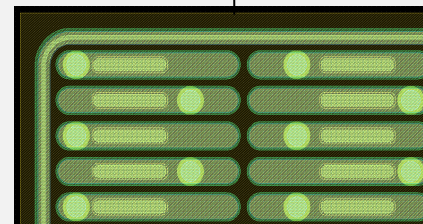
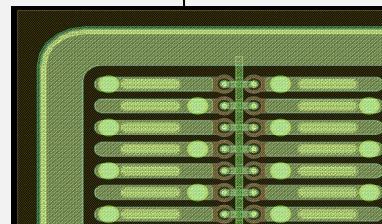
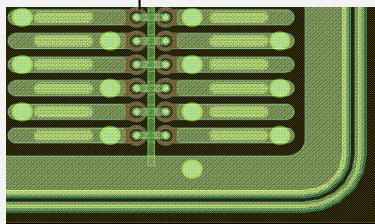
Thickness:

50 $\mu\text{m}$ , 100 $\mu\text{m}$ , 150 $\mu\text{m}$ .

UBM:

Platinum (Pt) or Nickel-Gold (NiAu).

Edge Design:



BR+GR	BR+No GR	GR+No BR	No BR+No GR
Inactive regions: 100 $\mu\text{m}$	78 $\mu\text{m}$	47 $\mu\text{m}$	47 $\mu\text{m}$

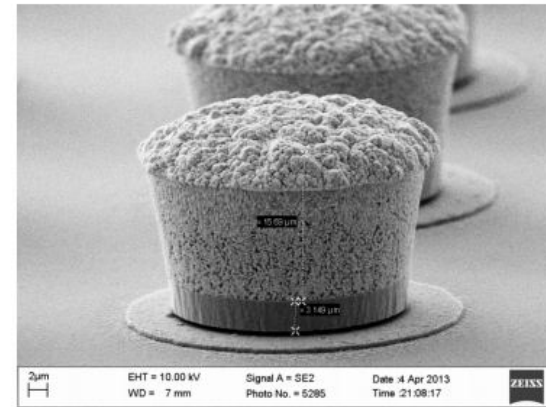
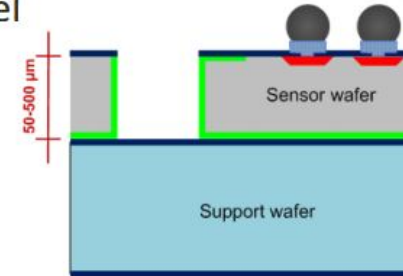
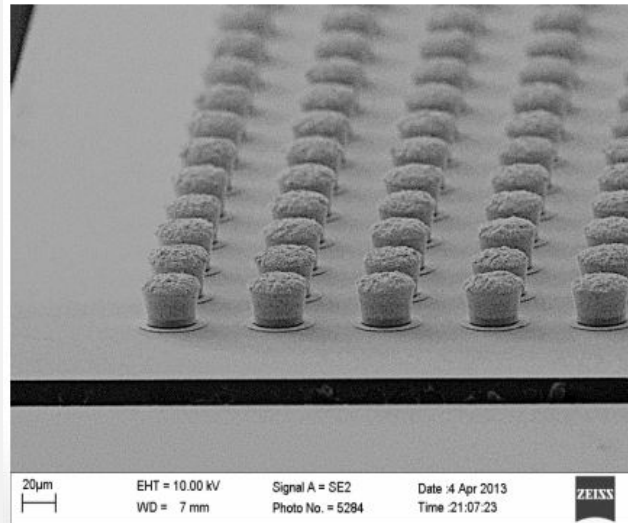


# New challenges on interconnections for Edgless pixel sensor

## Solder bumps on edgeless sensor

- Recently we have demonstrated a SnPb solder bump process on the edgeless sensors on a wafer level

*SnPb solders before reflow process*



April 9, 2013

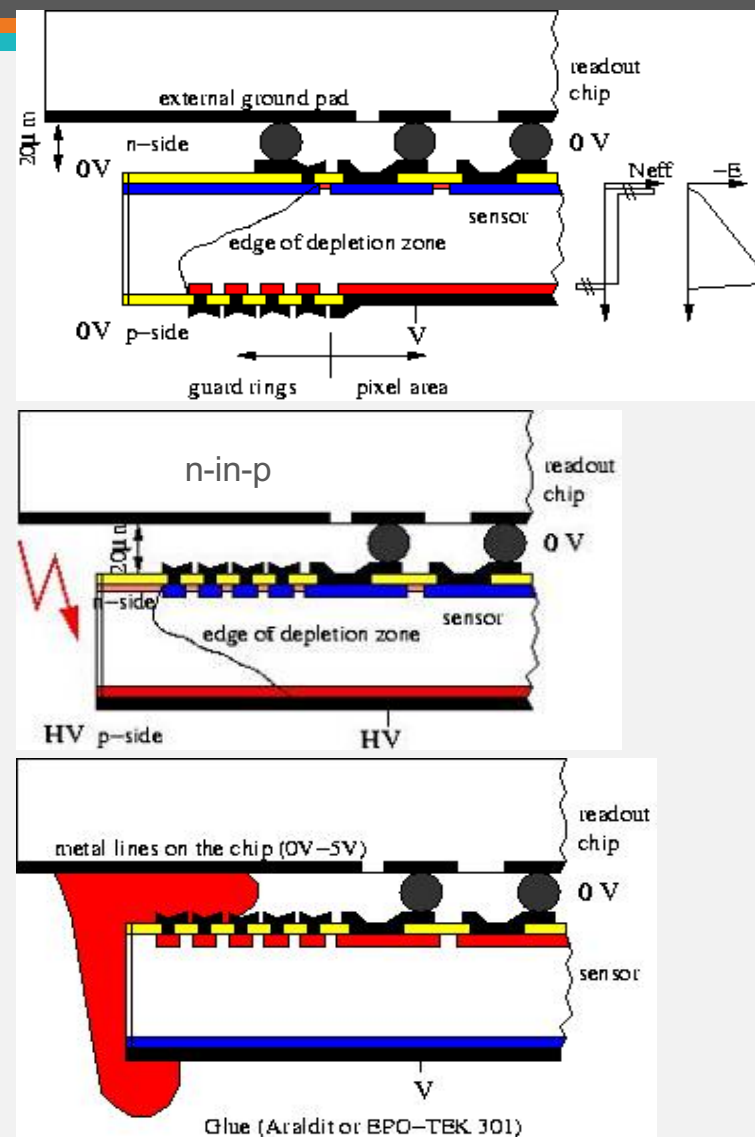
Juha.kallio@advacam.com

14

# R&D :Towards (n-in-p) Edgless Sensors for the future HL-LHC

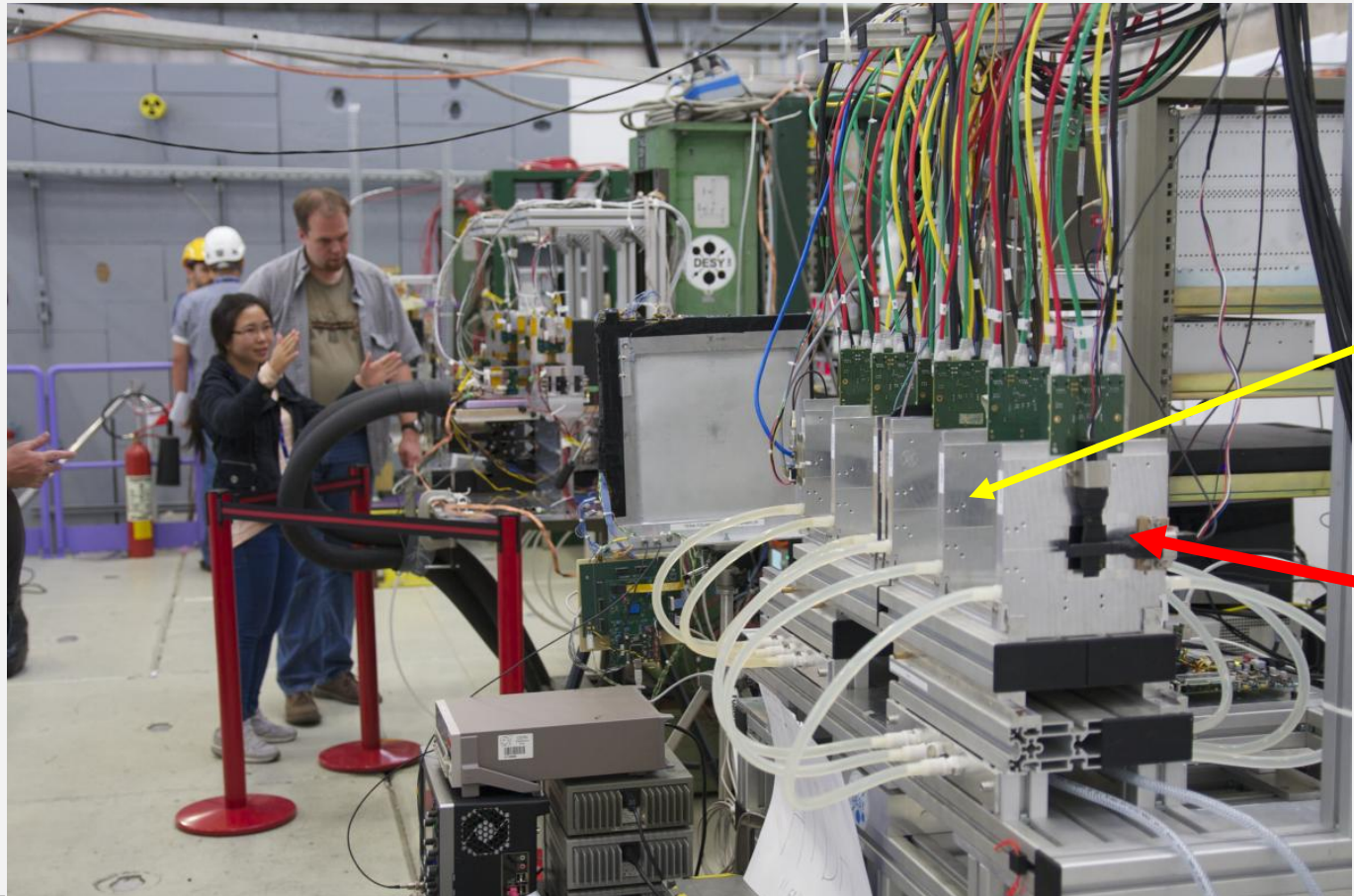
## ATLAS pixel detector uses n-in-n-sensors

- double sided processing (back side is structured)
- all sensor edges at ground
- most expensive part of the module
- *Exploring n-in-p sensors as alternative*
- Studies show radiation hardness
- single sided process ~ price benefit of factor 2-3
  - **Develop Active Edge Technology**
- Absence of guard rings on back side lead to risk of (destructive) sparking to the ROC



# Beam Test Telescope Tool

- A tool to study performance of our devices with high energy particles

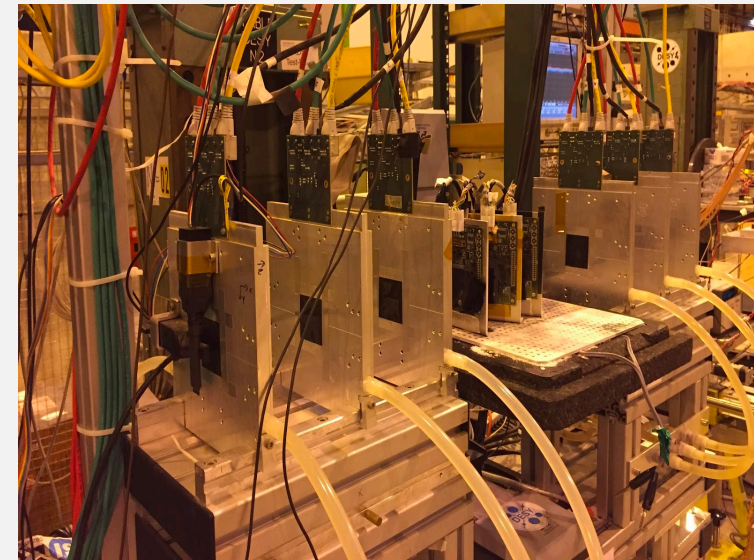
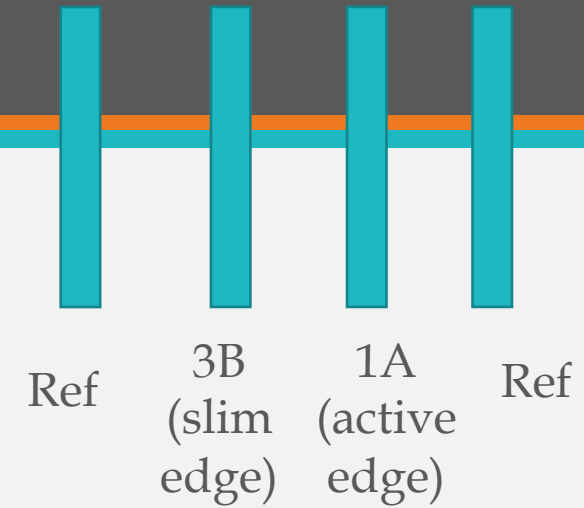




## Test Beam setup conditions

### Environment and purpose

- Beam: CERN SPS 120 GeV pion
- Telescopes: ACONITE (H6A) and AIDA (H6B)
  - Telescope planes with Mimosa26 sensors (1152x576), pixel pitch 18.4 $\mu$ m
- **Normal incidence (H6B #1)**
  - Dry ice cooling (approx. -40 to -50 °C\*)
  - Bias voltage points: 50 V, 80 V, 100 V, 120 V
- **Inclined track (H6B #2)**
  - DUT tilted at 27° (around y-axis)
  - Dry ice cooling (approx. -40 to -50 °C\*)
  - Bias voltage points: 50 V, 80 V, 100 V, 120 V



# Active edge and slim edge designs

- **Goal:**

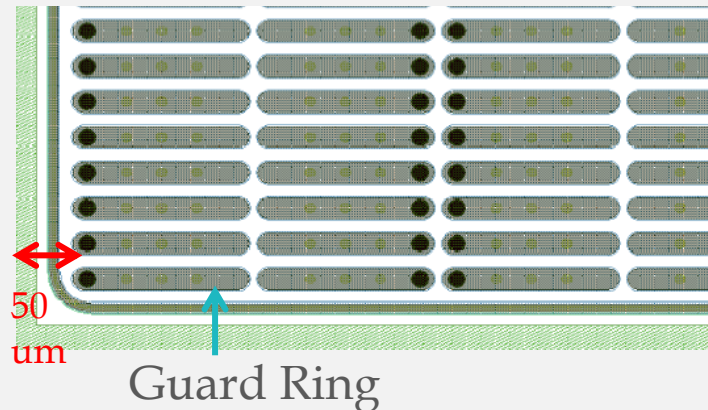
- Investigate the performance of the irradiated n-in-p planar pixel sensors with the active and slim edge design produced by ADVACAM;
- Study the overall and edge efficiency with normal and inclined tracks

- **Samples: active and slim edge design, 50  $\mu\text{m}$  thickness, low threshold tuning**

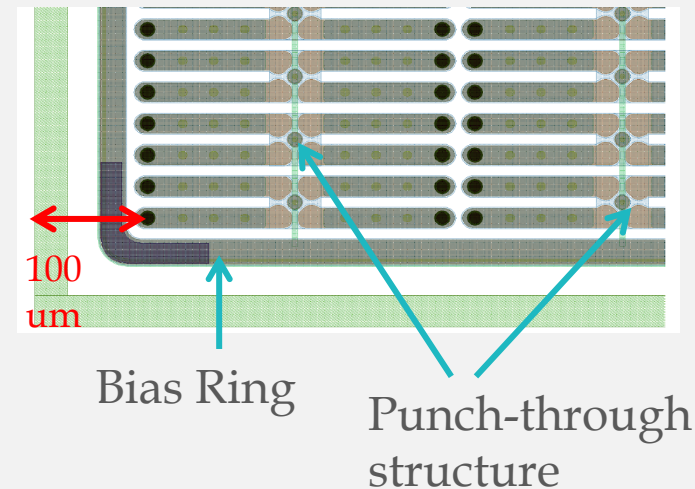
- Irradiation:  $\text{Irr.} = 1 \times 10^{15} \text{ n}_{\text{eq}}/\text{cm}^2$

Threshold : 700e, 8ToT4ke

Threshold : 520e-580e, 8ToT4ke

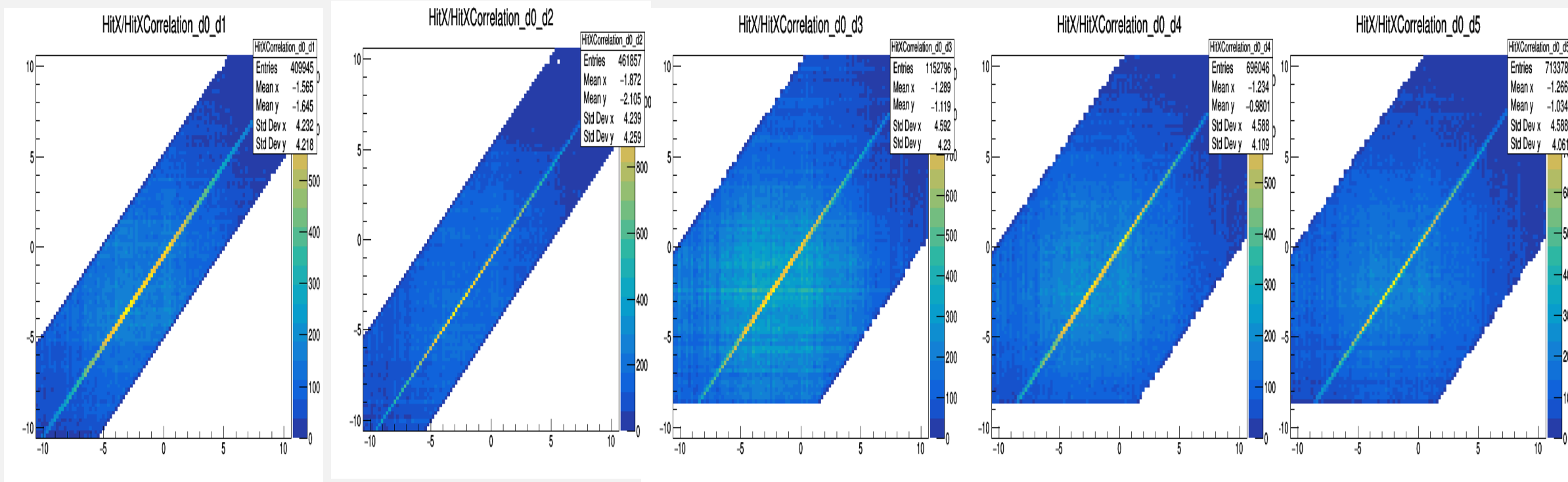


ADV-NP50-3-1A: active edge



ADV-NP50-3-3B: slim edge

# Correlations

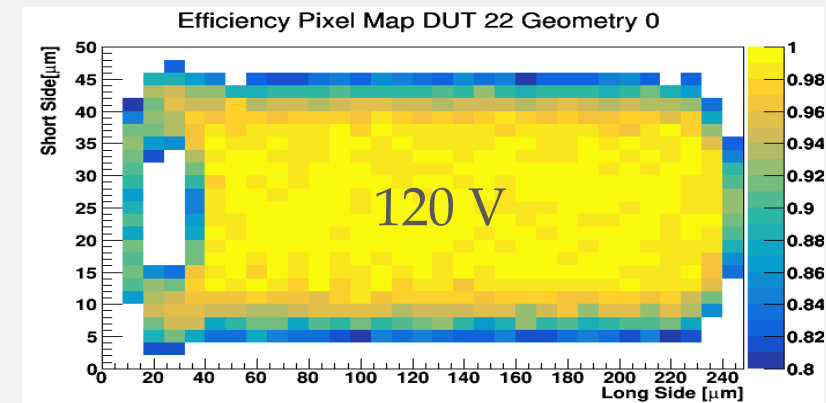
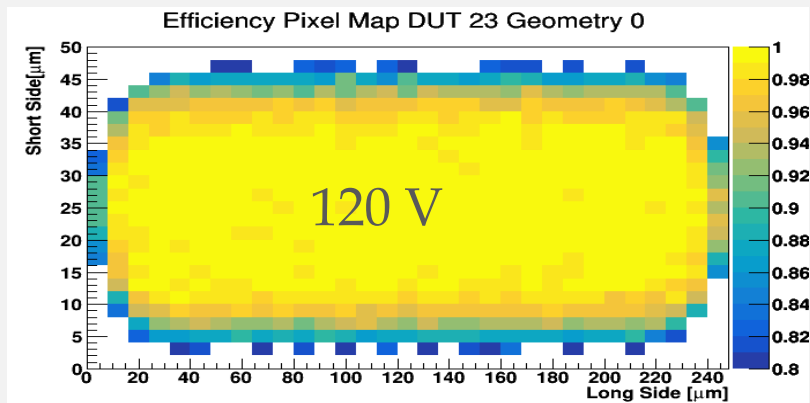
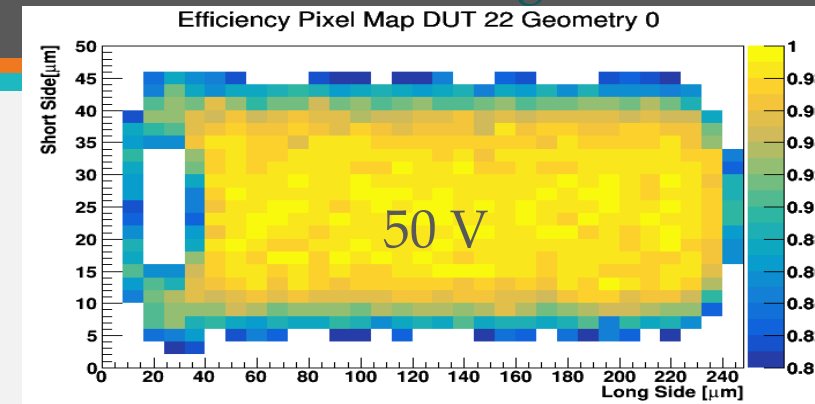
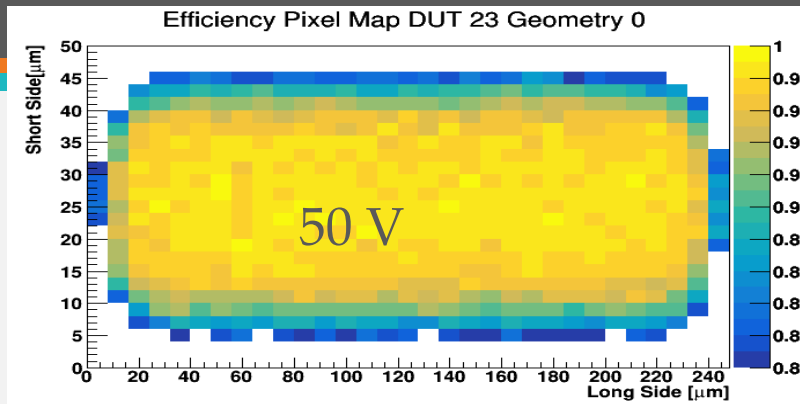


DUTs:  
normal

# In-pixel efficiency (normal tracks)

Active edge

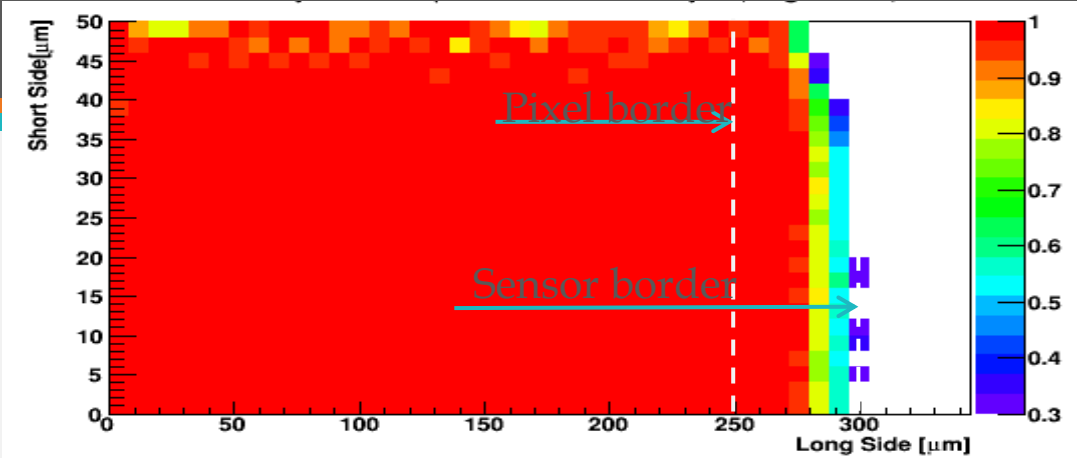
Slim edge



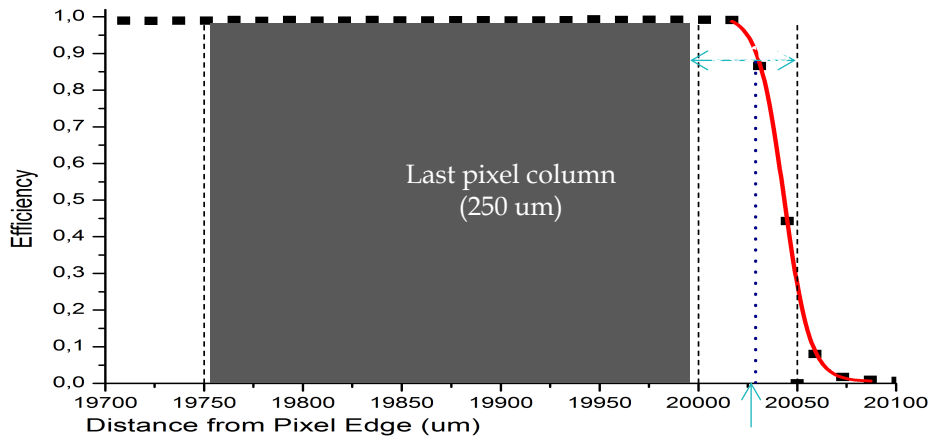
- Depletion region recovered with higher voltage in **central** pixel region
- Inefficiency in pixel sides and corners due to large charge sharing between cells for thin sensors

Testing the the performances at the borders to increase the overall surface efficiency

- Performances of pixels planaires under High energy Pions



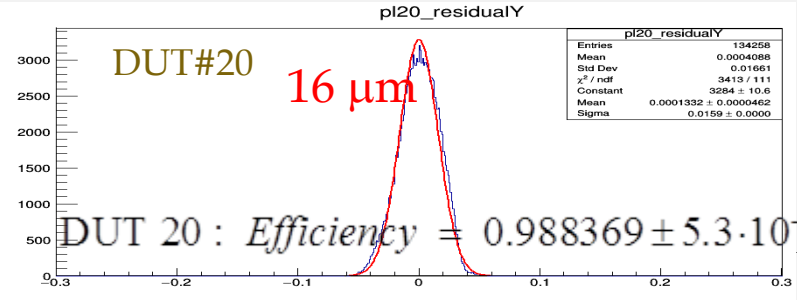
Active Edge 50 μm



AIDA HIGH RESOLUTION PIXEL TELESCOPE

D. Hohov, Phd Thesis  
University Paris Saclay 2019

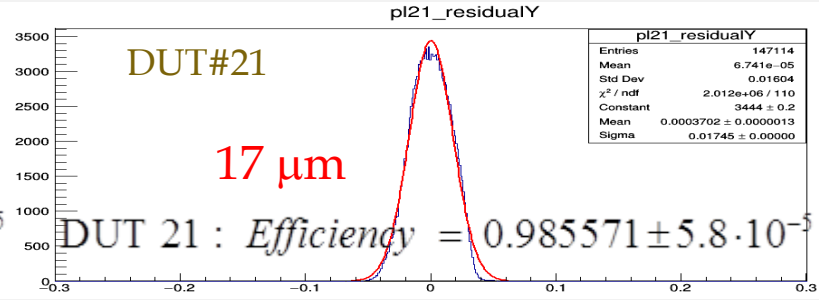
Distribution of residuals



DUT#20

16 μm

DUT 20 :  $Efficiency = 0.988369 \pm 5.3 \cdot 10^{-5}$

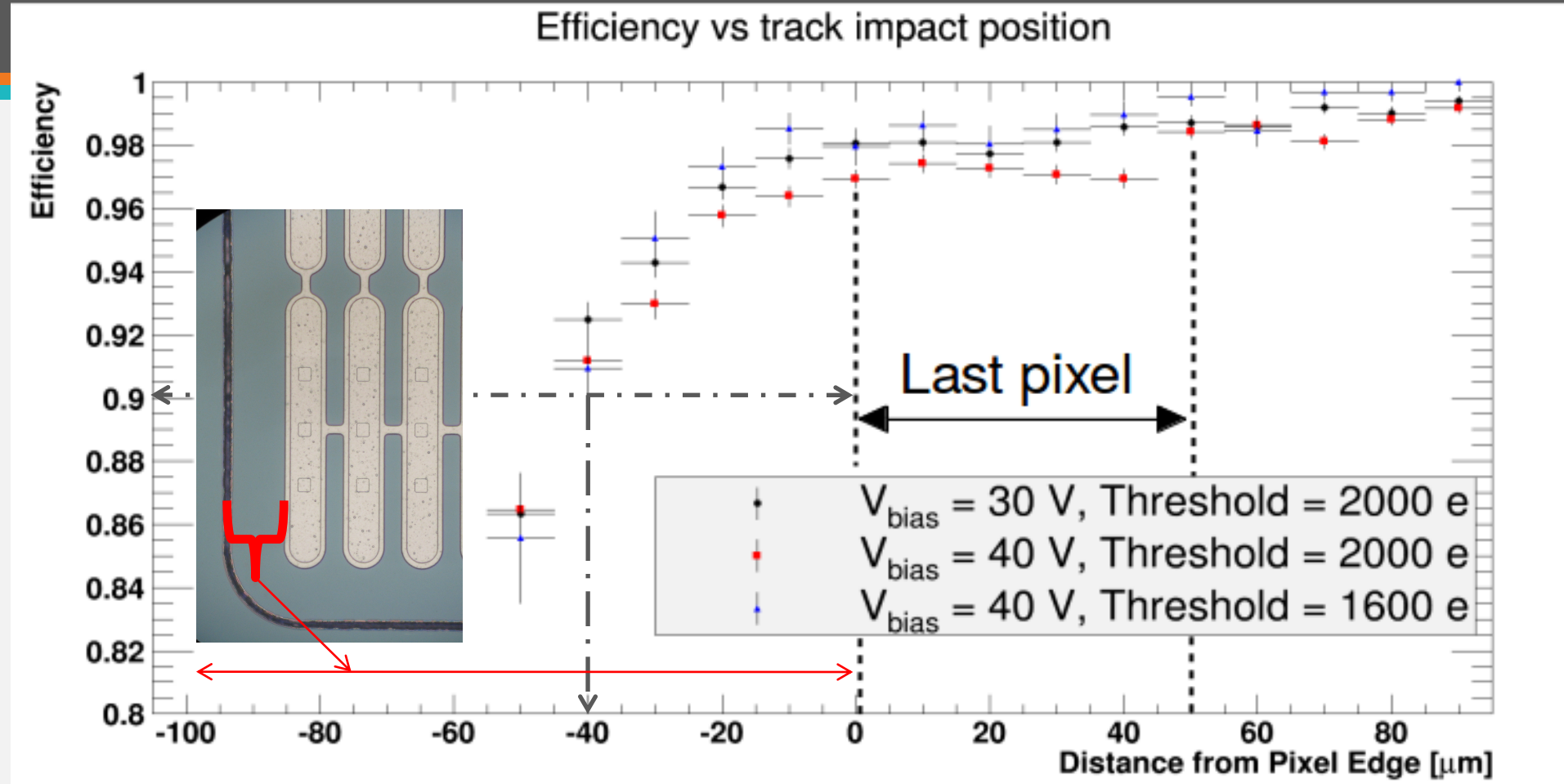


DUT#21

17 μm

DUT 21 :  $Efficiency = 0.985571 \pm 5.8 \cdot 10^{-5}$

# Hit-efficiency at the edge with High energy particles



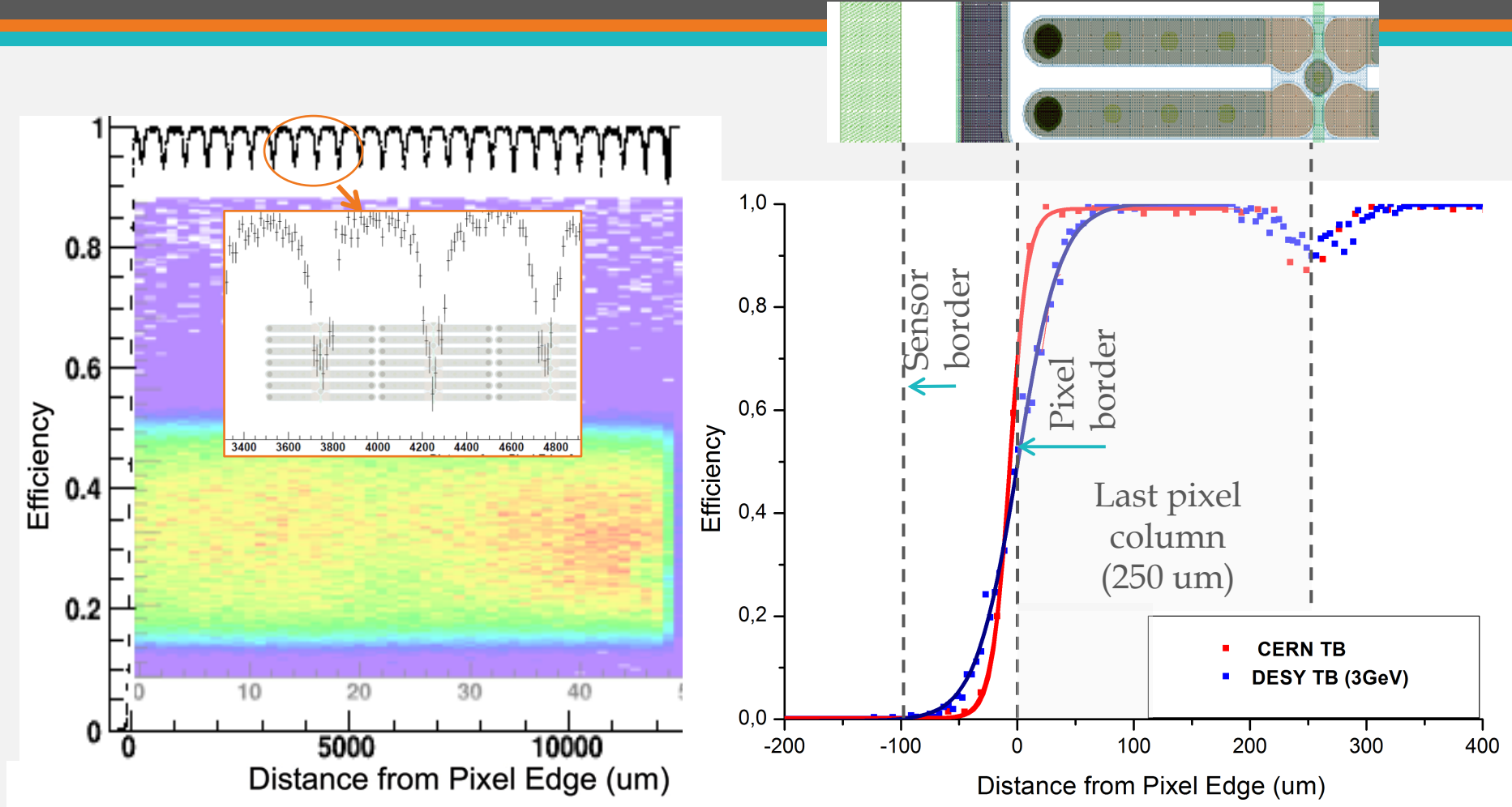
FRENCH-UKRAINIEN school for High Energy Physics and Medical Imaging

- Hit-efficiency above 90% up to 40  $\mu\text{m}$  away from the last pixel



# Efficiencies from beam tests

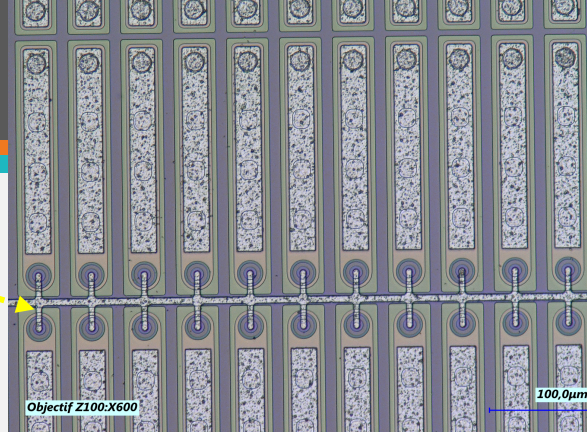
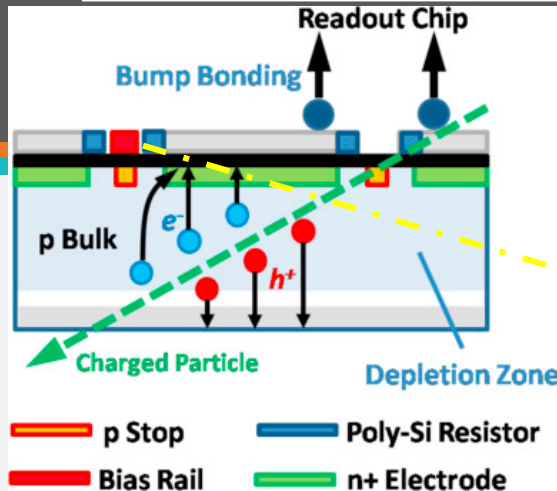
- *Efficiency vs track impact position DUT#22*



Analysis from Dmytro Hohov ATLAS@LAL

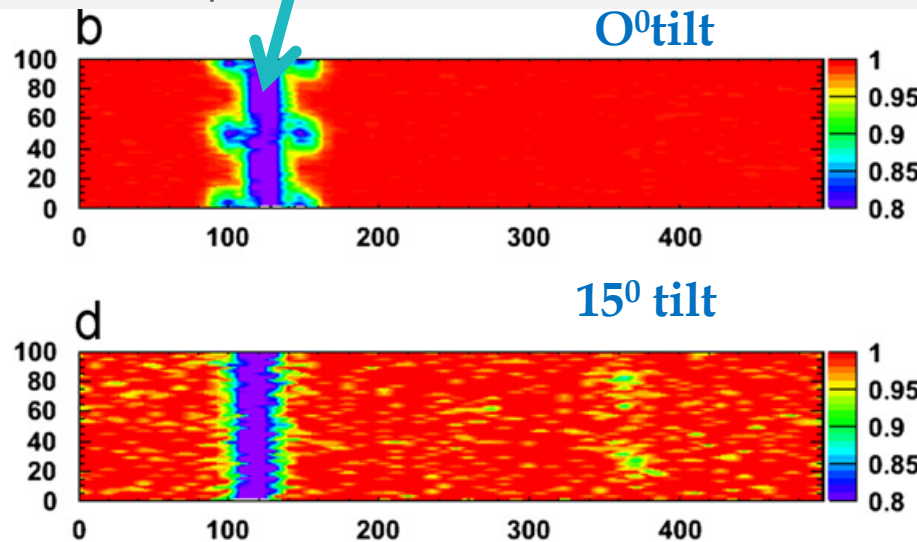
# Irradiated pixel behavior at high fluence

Beam-tests results



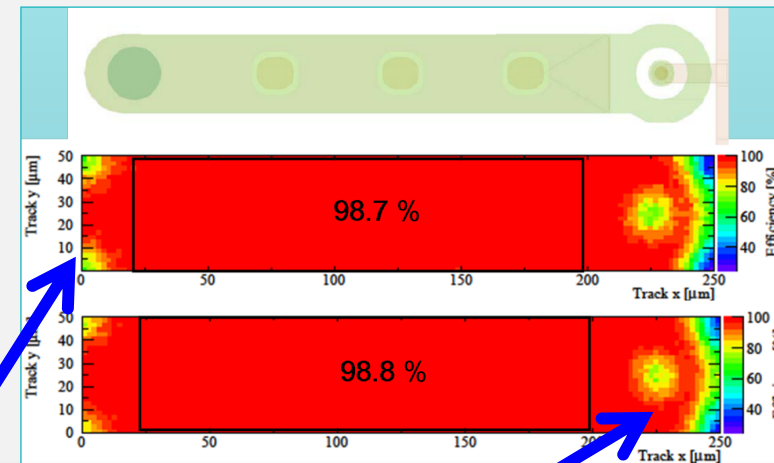
Bias Rail

Position of polarization rail



Pixel Matrix Hitmap at  $5 \cdot 10^{15}$  neq/cm<sup>2</sup> and Temperature -15°C.

Eudet 120 GeV pions/  $\Phi=4 \cdot 10^{15}$  n<sub>eq</sub> cm<sup>-2</sup>

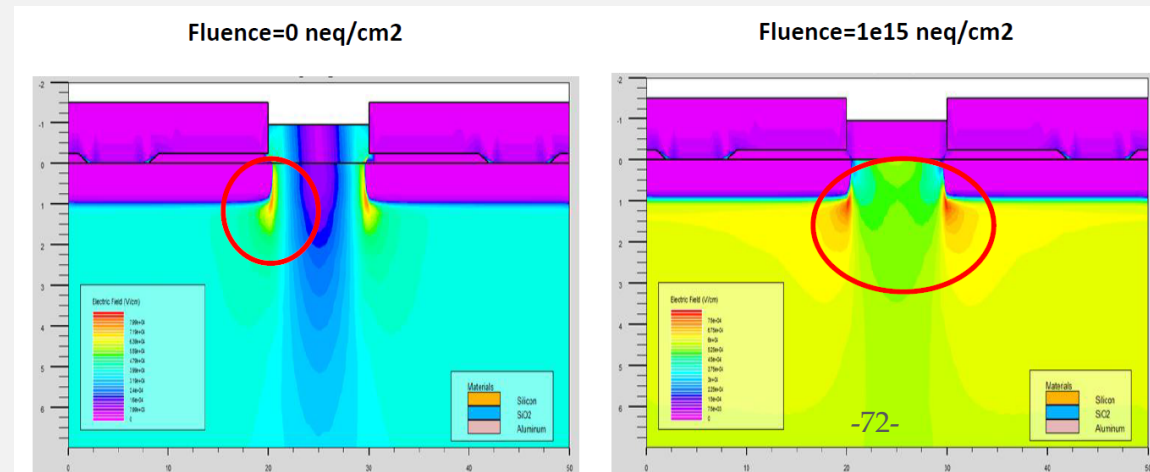
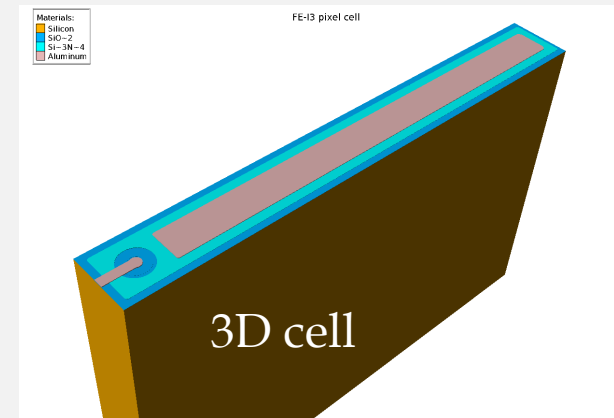


FRENCH-UKRAINIEN school for High Energy Physics and Medical Imaging Loss of Efficiency around Bias rails and implants

# An important TOOL for innovation : Simulation TeChnology Aided Design (TCAD) (Finite element Method)

- Aim: Develop simulations (TCAD input parameters) allowing to simulate performance of irradiated silicon sensors and performance predictions under various conditions (*sensor design and material, irradiation fluence and particle type, annealing, ...*).
  - Close collaboration with other ATLAS sensor simulation working groups (LAL member of RD50)
  - Challenge for irradiated sensors:
    - Correct implementation of bulk and surface damage by defect levels
    - Defect concentration is function of fluence, particle type, material, annealing, .... !
  - Validate the radiation damage model to be implemented in TCAD packages to be able to simulate the performance of complex silicon devices after hadron irradiation. The simulation output is increasingly accurate in term of IV, CV, CCE inter-electrode resistance and capacitance, break-down voltage prediction etc.

Evaluating different pixel layout designs and production parameters (p-stop, p-spray)



## Finite Element Simulation

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- **Our Problem: Solution of Laplace Equation and Continuity equations in regions**

**Physics models:** Works by modelling electrostatic potential (Poisson's equation) and carrier continuity (drift-diffusion, dd, mainly)

Poisson  $\epsilon_s \nabla \cdot \underline{E} = -\epsilon_s \nabla^2 \psi = q(p - n + N)$

Electron continuity  $\frac{\partial n}{\partial t} = \frac{1}{q} \nabla \cdot \underline{J}_n + (G - R)$   $\underline{J}_p = q\mu_n \underline{E} - qD_p \nabla p$

Hole continuity  $\frac{\partial p}{\partial t} = \frac{-1}{q} \nabla \cdot \underline{J}_p + (G - R)$   $\underline{J}_n = q\mu_n \underline{E} + qD_n \nabla n$

See Fichtner, Rose, Bank, "Semiconductor Device Simulation", IEEE Trans. Electron Devices 30 (9), pp1018, 1983

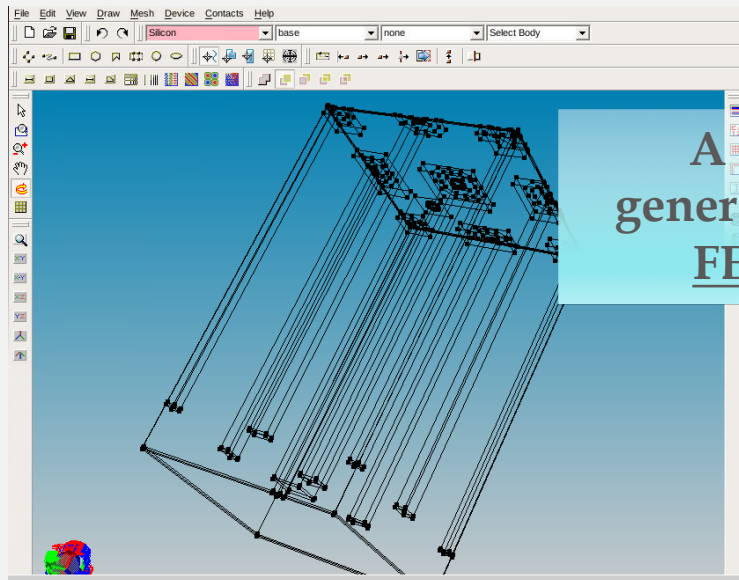
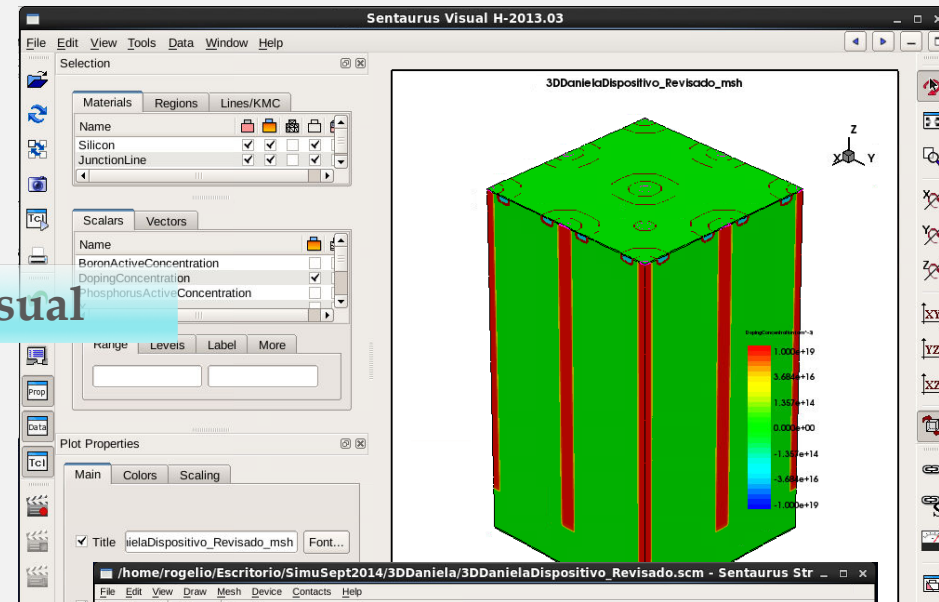
Different versions of physics models available

- Different models of mobility, bandgap...
- Generation and recombination rates may include avalanche effects, charge generation by high-energy particles...

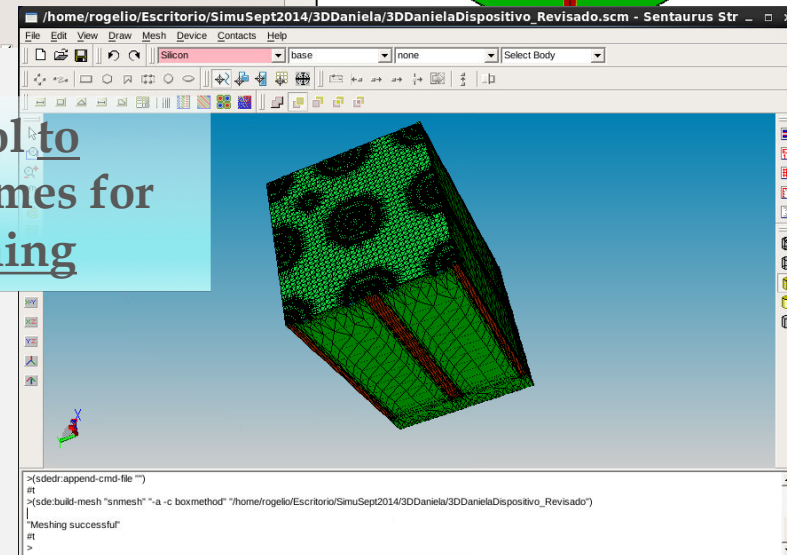
# Finite Element Simulation

- Basic tools in a Finite Element simulator
  - Automatic Mesher Creator tool  
(Sentaurus Structure Editor, SSE)
  - Visualization tool  
(SVisual)
  - Solver tool  
Sentaurus Device

SVisual



A CAD tool to  
generate volumes for  
FEM meshing

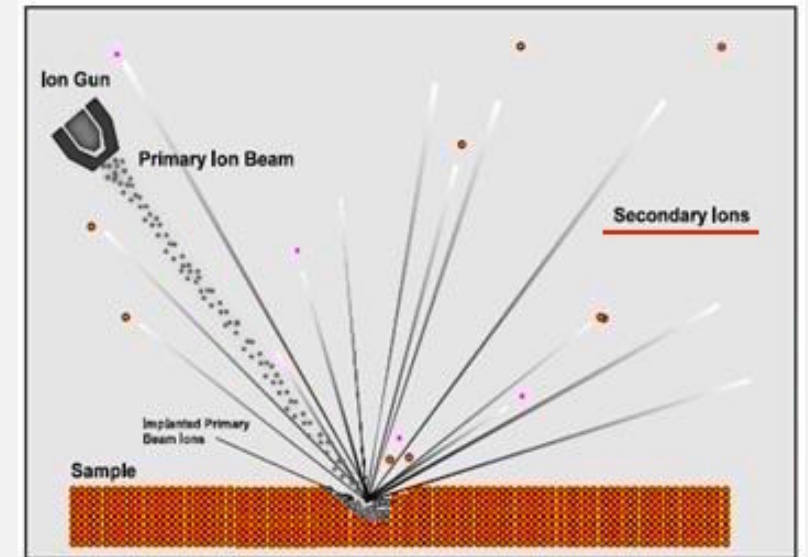




# *A Tool to help understanding Pixel devices :*

## *Secondary Ion Mass Spectrometry (SIMS)*

- SCANNING of intrinsic material : doping profile.
- secondary ions ejected from a sample surface when bombarded by a primary beam allows you to find out the atomic structure of your sample



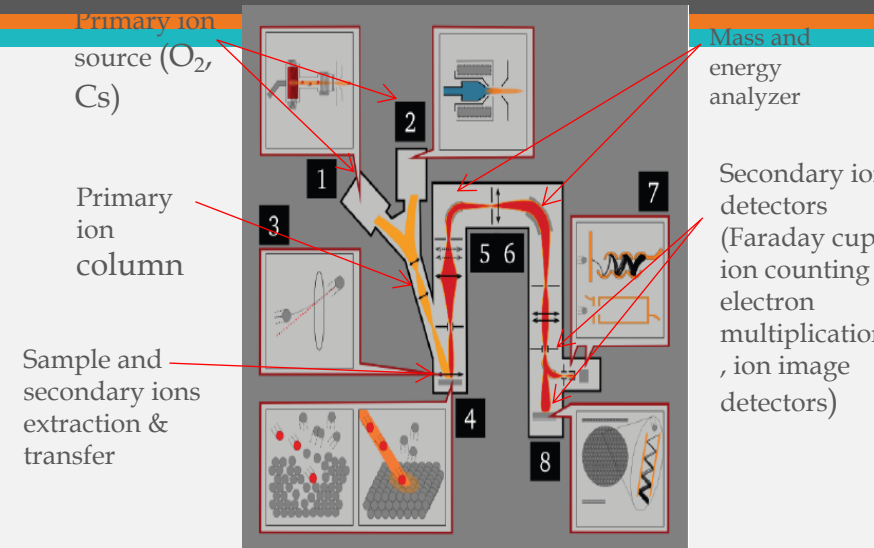
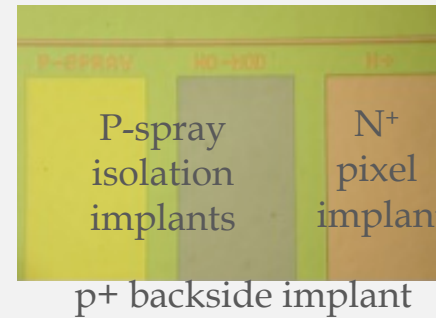
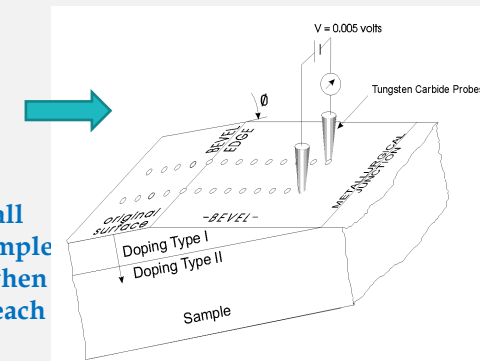


# Design improvements Tools: synopsis 3D & doping profile measurements

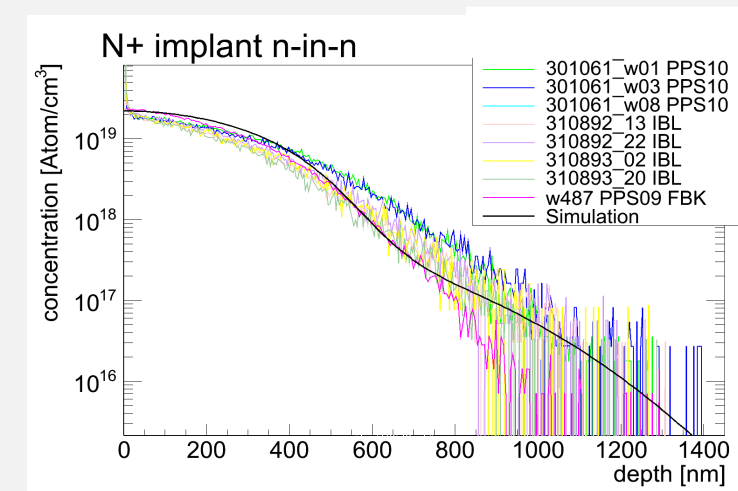
- To overcome the problem of non-accessibility of Technology parameters :
- charge carrier distribution, doping profiles, density of defects and Impurities etc...
- TCAD simulation is absolutely needed
- Tool for studying relationship: semiconductor physics - device characteristics
- Process simulator: oxidation, ion implantation, diffusion, etching etc....
- Device simulator: electrical parameters by solving eq.: Poisson, current density , continuity etc..

- Experimental methods used for TCAD calibration
- Secondary Ion Mass Spectrometry (SIMS) – atomic doping profile
- Spreading Resistance Profiling (SRP) – charge carriers doping profile

SRP : It relies on stepping a pair of small probes across the bevel surface of a sample and the measurement of a resistance when 5 mV are applied across the probes at each step



- Optimization of implantation & annealing parameters:
  - energy and dose, oxide thickness, annealing time and temperature



## 3D doping Profile Measurement : New Method

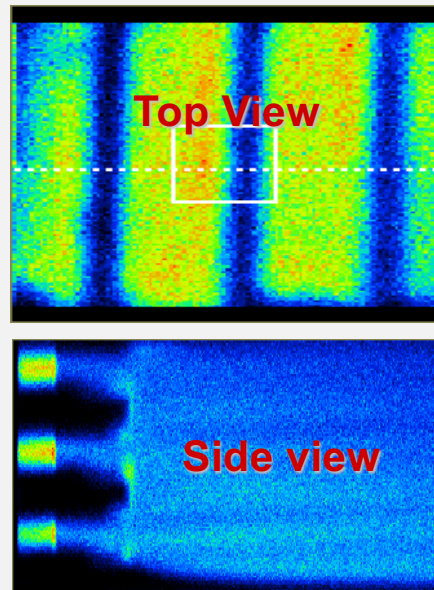
- **Standard SIMS** very helpful tool for: process control, in particular diffusion, identification of contaminants and failure in the fabrication.
- Results for **1D doping profiles** of dopant are satisfactory.
- Now, we want to study the doping profile at the pixel level of an **Advacam Active Edge Detectors**.
  - Analyzing small region of interest like the pixel region and the active edge region, need higher lateral resolution technique.

we need to move to another method called SIMS Imaging.

# SIMS Imaging Method

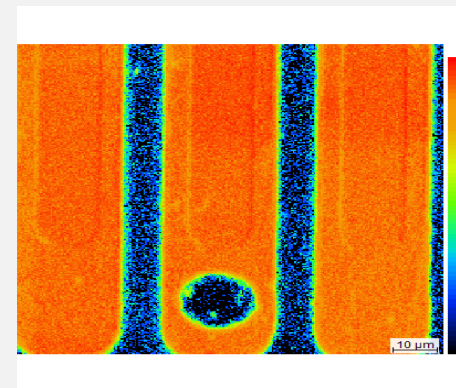
## •What is SIMS imaging?

- Allow a scan for the samples surface and depth.
- Depth profiling and imaging can be combined to yield very powerful **three-dimensional dopant maps**



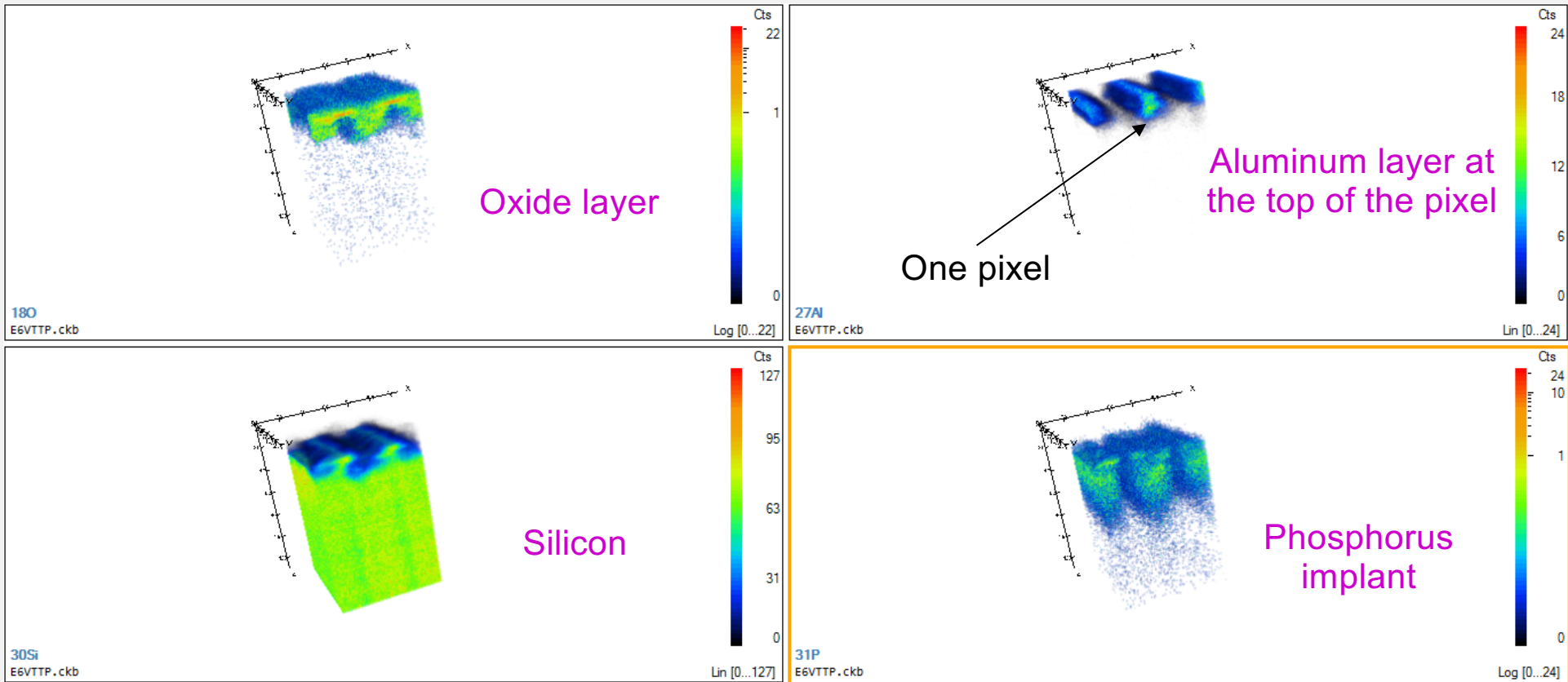
## Why SIMS imaging?

- ✓ Can achieve lateral resolutions up to 5  $\mu\text{m}$ .
- ✓ High surface sensitivity at ppb level can be reached.
- ✓ Sample preparation is rather simple.
- ✓ Equivalent measuring time with standard 1D SIMS.



## Results: Phosphorus Implant 3D Doping Map

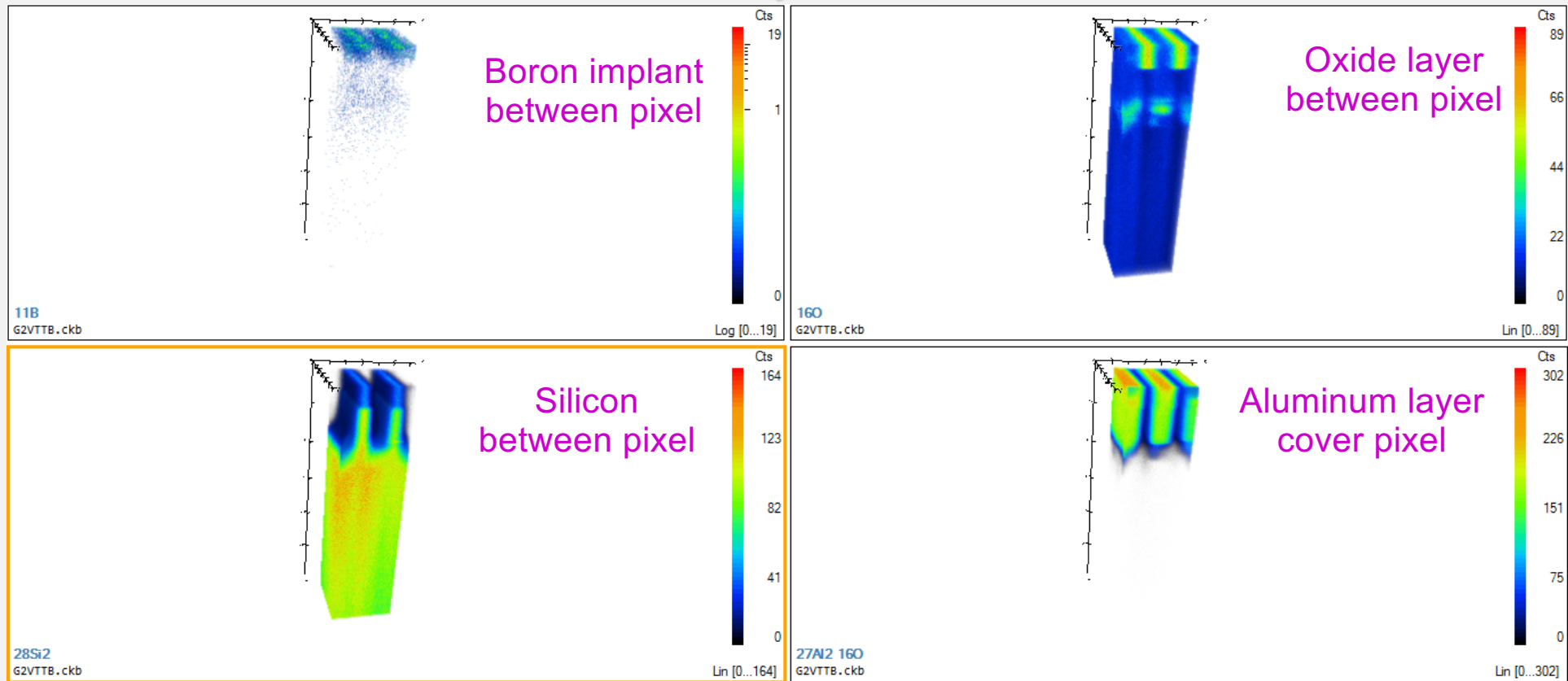
- Analyzed area: Center Pixel region.



# Results: p-spray Boron Implant between the pixel

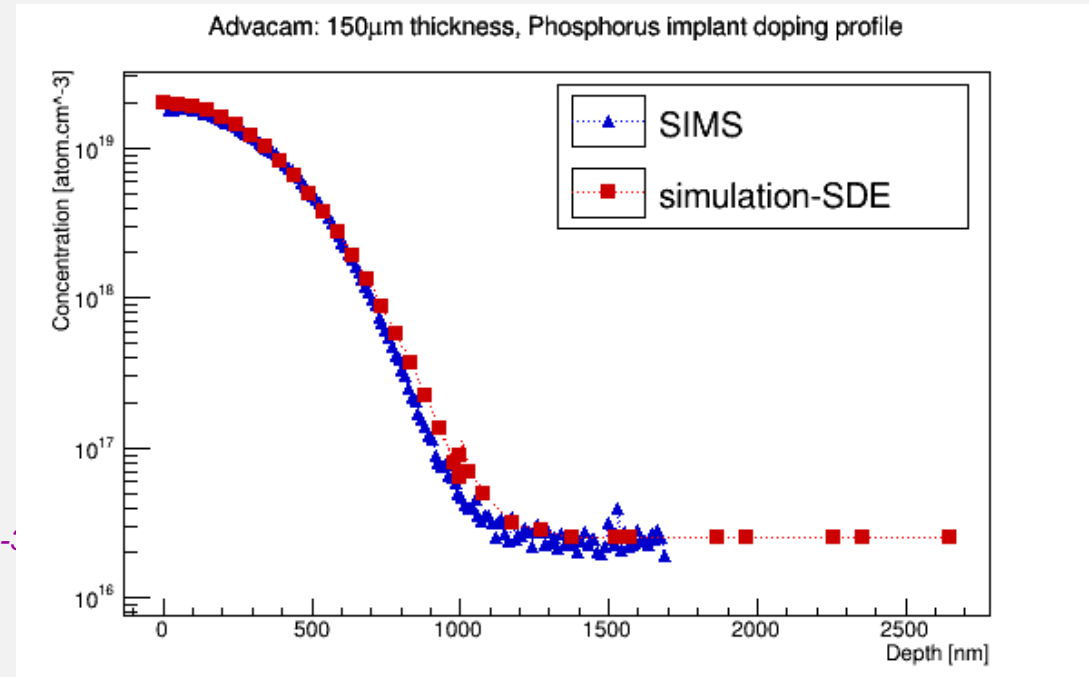
- Analyzed area: between pixel region.

## Preliminary results



# Results: Phosphorus Implant Doping profile

- Analyzed area: Center Pixel region: SIMS-Simulation comparison.
- Wafer 150-2: 150  $\mu\text{m}$ , BR+GR design
- **SDE** Synopsys simulation.
- Simulation parameter:
  - Analytical doping profile
  - Peak Value concentration  $2\text{E}19$
  - Depth of Implant  $1.5\text{ }\mu\text{m}$ .
- The implant extends  $2\text{ }\mu\text{m}$  in the silicon substrate.
- Implant peak concentration  $1\text{E}19\text{ atom.cm}^{-3}$
- Oxide layer has not been simulated.

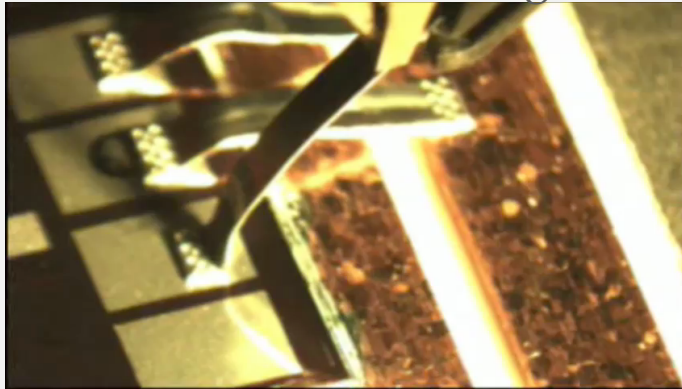




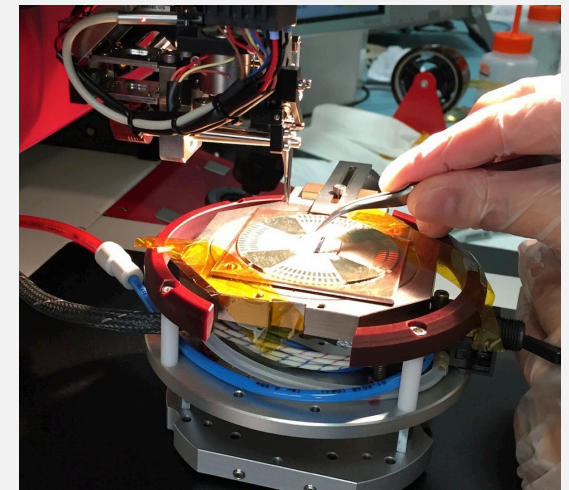
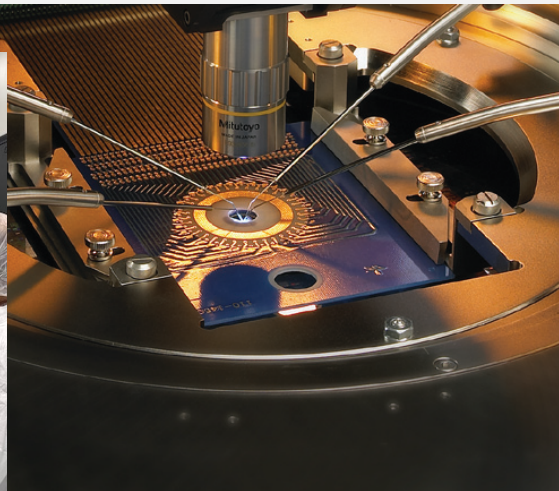
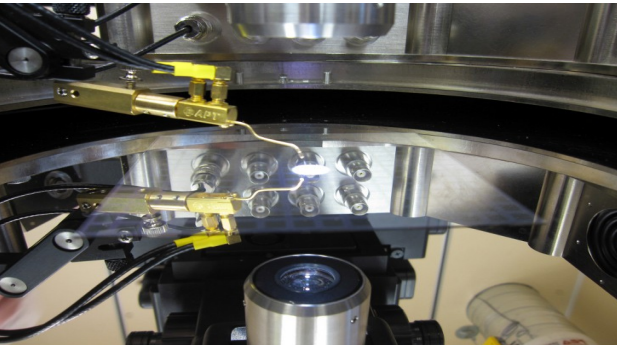
# State of the art Infrastructure



Wire Bonding machine



Clean room

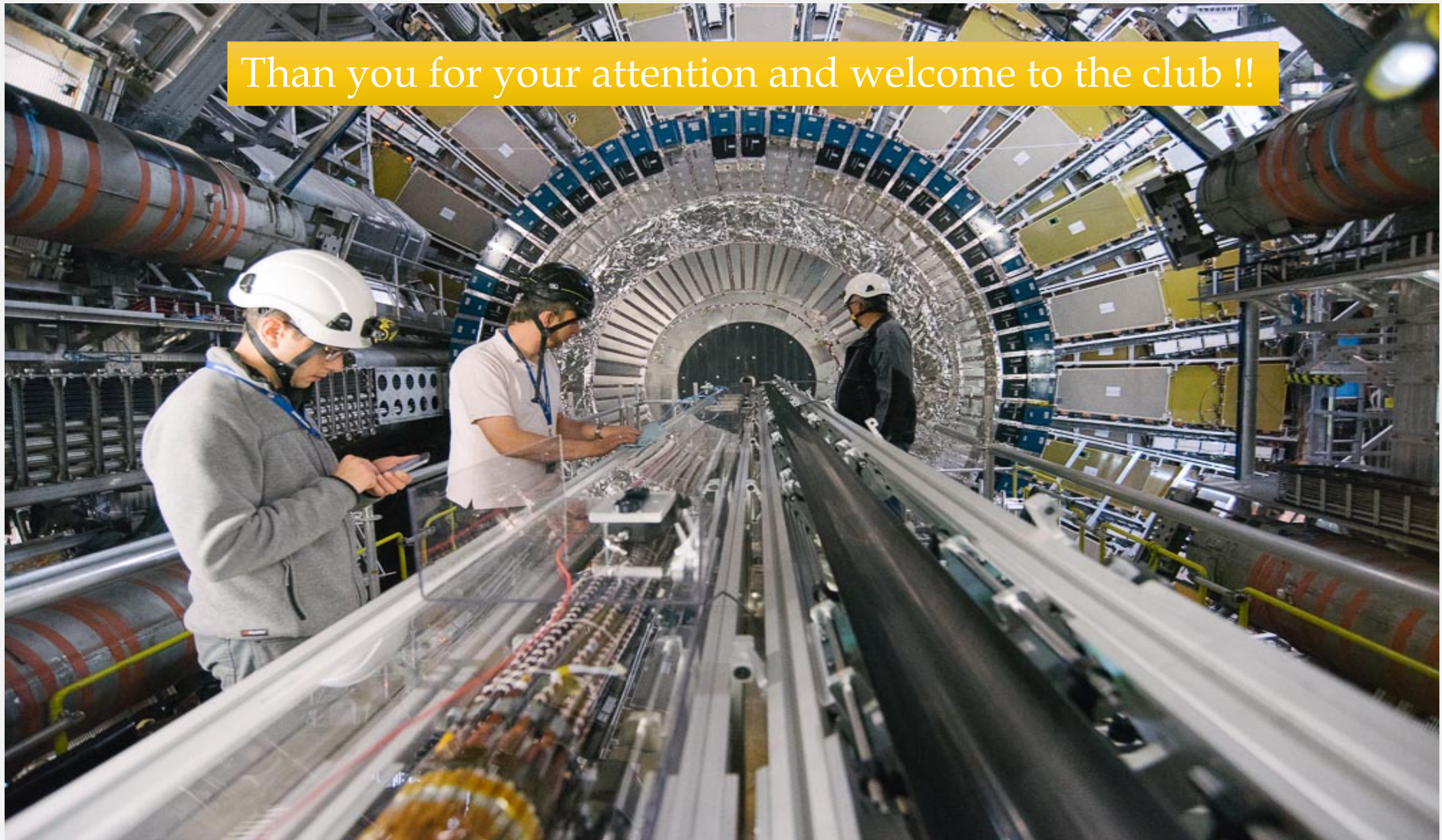


# conclusions

- R&D on planar pixel sensors is driven by ATLAS High luminosity upgrade
- The goal is to optimize the design of the pixel sensors to cope
  - High radiation levels : radio-tolerant materials
  - High level of occupancy, favours high granular sensors
  - Develop innovative approach for device bulk type
  - Increase full sensor sensitivity , go to edgeless technology
  - Predict, improve new device concepts by TCAD simulations
  - Still room for improvement and innovations ....



Than you for your attention and welcome to the club !!

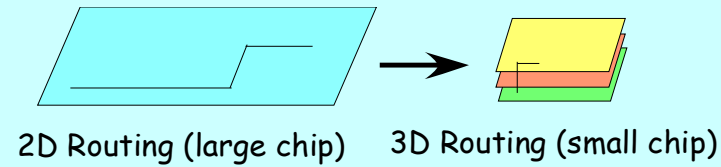


*The end*

*The ideal case : Physicist dream*

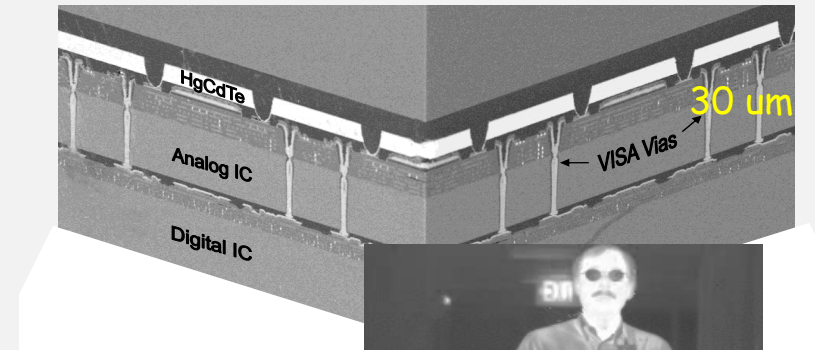
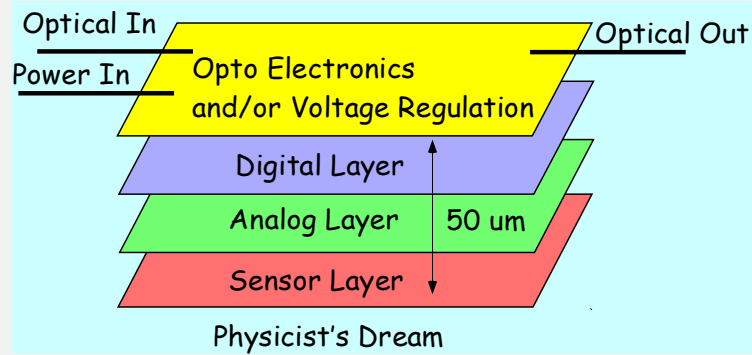
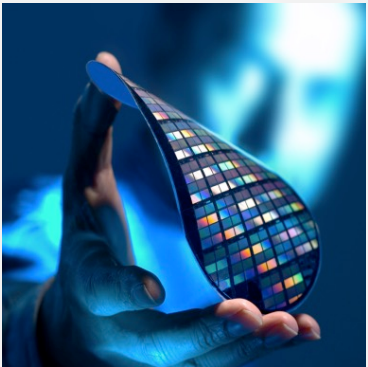


# Vertical Scale Integration (3D)



idea: implement thin layers of silicon, each with independant specific function.

Future



Advantages: one object !! Monolithic!!  
Reduction of  $R, L, C$   
Lower  $X_0$ , better irradiation tolerance  
low consumption, low capacity  
Increase fonctionnalities (intelligence in chip)

Through silicon vias

